# Imperial College London MSci EXAMINATION May 2014 

This paper is also taken for the relevant Examination for the Associateship

## GENERAL RELATIVITY

For 4th-Year Physics Students
Monday $19^{\text {th }}$ May 2014: 14:00 to 16:00

The paper consists of two sections: $A$ and $B$
Section A contains one question [40 marks total].
Section B contains four questions [30 marks each].
Candidates are required to:
Answer ALL parts of Section A and TWO QUESTIONS from Section B.
Marks shown on this paper are indicative of those the Examiners anticipate assigning.

## General Instructions

Complete the front cover of each of the 3 answer books provided.
If an electronic calculator is used, write its serial number at the top of the front cover of each answer book.

## USE ONE ANSWER BOOK FOR EACH QUESTION.

Enter the number of each question attempted in the box on the front cover of its corresponding answer book.

Hand in 3 answer books even if they have not all been used.
You are reminded that Examiners attach great importance to legibility, accuracy and clarity of expression.

## Conventions:

We use conventions as in lectures. In particular we take (-,+,+,+) signature.

## You may find the following formulae useful:

The Christoffel symbol is defined as,

$$
\Gamma_{\alpha \beta}^{\mu} \equiv \frac{1}{2} g^{\mu v}\left(\partial_{\alpha} g_{v \beta}+\partial_{\beta} g_{\alpha v}-\partial_{\nu} g_{\alpha \beta}\right)
$$

The covariant derivative of a vector field is,

$$
\nabla_{\mu} v^{v} \equiv \partial_{\mu} v^{v}+\Gamma^{v}{ }_{\mu \alpha} v^{\alpha}
$$

and for a covector field is,

$$
\nabla_{\mu} w_{\nu} \equiv \partial_{\mu} w_{\nu}-\Gamma^{\alpha}{ }_{\mu \nu} w_{\alpha}
$$

For a Lagrangian of a curve $x^{\mu}(\lambda)$ of the form,

$$
L=\int d \lambda \mathcal{L}\left(x^{\mu}, \frac{d x^{\mu}}{d \lambda}\right)
$$

the Euler-Lagrange equations are,

$$
\frac{d}{d \lambda}\left(\frac{\partial \mathcal{L}}{\partial\left(\frac{d x^{\mu}}{d \lambda}\right)}\right)=\frac{\partial \mathcal{L}}{\partial x^{\mu}}
$$

## Section A

Answer all of section A.

## SECTION A

1. This question concerns the covariant derivative.
(i) State how the components of a $(1,0)$ tensor $v^{\mu}$ and a $(0,1)$ tensor $w_{\mu}$ transform under a coordinate transformation $x \rightarrow x^{\prime}$.
[8 marks]
(ii) Use your previous answer to show that $v^{\mu} w_{\mu}$ transforms as a scalar under a coordinate transformation $x \rightarrow x^{\prime}$.
[8 marks]
(iii) Under a coordinate transformation the Christoffel symbol transforms as;

$$
\Gamma_{\alpha^{\prime} \beta^{\prime}}^{\prime \mu^{\prime}}=\Gamma_{\alpha \beta}^{\mu} \frac{\partial x^{\prime \mu^{\prime}}}{\partial x^{\mu}} \frac{\partial x^{\alpha}}{\partial x^{\prime \alpha^{\prime}}} \frac{\partial x^{\beta}}{\partial x^{\prime \beta^{\prime}}}-\left(\frac{\partial^{2} x^{\prime \mu^{\prime}}}{\partial x^{\alpha} \partial x^{\beta}}\right) \frac{\partial x^{\alpha}}{\partial x^{\prime \alpha^{\prime}}} \frac{\partial x^{\beta}}{\partial x^{\prime \beta^{\prime}}} .
$$

Show that the Christoffel symbol does not transform as a tensor.
[2 marks]
(iv) Show that $\partial_{\mu} w_{v}$, the partial derivative of a covector field $w_{\mu}$, does not transform as a tensor.
(v) Starting from the relations,

$$
\delta_{v}^{\mu}=\frac{\partial x^{\mu}}{\partial x^{v}}=\frac{\partial x^{\mu}}{\partial x^{\prime v^{\prime}}} \frac{\partial x^{\prime v^{\prime}}}{\partial x^{v}}
$$

take an appropriate partial derivative of this to derive,

$$
\frac{\partial x^{\prime \alpha}}{\partial x^{\alpha}} \frac{\partial x^{\prime \beta}}{\partial x^{\beta}} \frac{\partial^{2} x^{\mu}}{\partial x^{\prime \alpha} \partial x^{\prime \beta^{\prime}}}=-\frac{\partial x^{\mu}}{\partial x^{\prime \mu^{\prime}}} \frac{\partial^{2} x^{\prime \mu \mu^{\prime}}}{\partial x^{\alpha} \partial x^{\beta}} .
$$

[6 marks]
(vi) Show that the covariant derivative of a covector field $w_{\mu}$, defined as $\nabla_{\mu} w_{v}=$ $\partial_{\mu} w_{v}-\Gamma^{\alpha}{ }_{\mu \nu} W_{\alpha}$, does transform as a tensor.

## Section B

Answer 2 out of the 4 questions in the following section.

## SECTION B

2. This question concerns the Newtonian spacetime, which we write using coordinates $x^{\mu}=\left(t, x^{i}\right)$ with $i=1,2,3$ as,

$$
d s^{2}=\left(\eta_{\mu \nu}-2 \epsilon \Phi\left(x^{i}\right) \delta_{\mu \nu}\right) d x^{\mu} d x^{\nu}
$$

where $\epsilon \Phi$ is the Newtonian potential, and we are interested in the Newtonian limit $\epsilon \rightarrow 0$. We take $\Phi$ to be static, and hence only a function of the $x^{i}$.
(i) State the stress tensor for a perfect fluid in a general spacetime in terms of its energy density $\rho$, pressure $P$ and local 4 -velocity $u^{\mu}$ (where $u^{\mu} u_{\mu}=-1$ ). What conservation equation does the stress tensor obey?
(ii) In the limit $\epsilon \rightarrow 0$ the Ricci tensor to leading order $O(\epsilon)$ is;

$$
\begin{aligned}
R_{t t} & =\epsilon \delta_{i j} \partial_{i} \partial_{j} \Phi \\
R_{t i} & =0 \\
R_{i j} & =\epsilon \delta_{i j}\left(\delta_{k l} \partial_{k} \partial_{l} \Phi\right)
\end{aligned}
$$

Use this to compute the components of the stress tensor that satisfies the Einstein equations for this spacetime. Show that this is the stress tensor for a dust fluid (ie. fluid with zero pressure), and determine the 4 -velocity and energy density of this dust in terms of the Newtonian potential $\epsilon \Phi$.
[12 marks]
(iii) By calculation, show that to leading order in $\epsilon$,

$$
\Gamma_{t t}^{i}=\epsilon \partial_{i} \Phi
$$

Using this, show that a non-accelerated particle with proper time $\tau$ that is slowly moving obeys (to leading order),

$$
\frac{d^{2} x^{i}}{d \tau^{2}}=-\partial_{i}(\epsilon \Phi)
$$

[10 marks]
(iv) Use these answers to briefly explain how Newton's force law of gravity arises in General Relativity.
[2 marks]
[Total 30 marks]
3. This question concerns the Schwarzschild metric, which we write using coordinates $x^{\mu}=(t, r, \theta, \phi)$ as,

$$
d s^{2}=-\left(1-\frac{2 G M}{r}\right) d t^{2}+\left(1-\frac{2 G M}{r}\right)^{-1} d r^{2}+r^{2}\left(d \theta^{2}+\sin ^{2} \theta d \phi^{2}\right)
$$

for a mass $M$, with $G$ the Newton constant.
(i) Consider a timelike geodesic $x^{\mu}(\tau)=(T(\tau), R(\tau), \Theta(\tau), \Phi(\tau))$ in the Schwarzschild metric where $\tau$ is proper time. Write a Lagrangian that we may vary to determine the geodesic. Deduce the Euler-Lagrange equations for $\Theta$ and $\Phi$ and show these are consistent with a geodesic that lies in the plane $\theta=\pi / 2$. We now restrict our attention to such geodesics. Show then that,

$$
R^{2} \frac{d \Phi}{d \tau}=J
$$

where $J$ is a constant.
(ii) Further deduce the equations that govern $T$ and $R$. Show that,

$$
\left(1-\frac{2 G M}{r}\right) \frac{d T}{d \tau}=k
$$

where $k$ is a constant. Hence show the equation governing the radial motion in the plane $\theta=\pi / 2$ looks like that of one dimensional motion for a unit mass particle in a potential $V(R)$ with constant energy $E$ so,

$$
E=\frac{1}{2}\left(\frac{d R}{d \tau}\right)^{2}+V(R), \quad V(R)=-\frac{G M}{R}+\frac{J^{2}}{2 R^{2}}+\frac{\alpha J^{2}}{R^{3}}
$$

where $\alpha$ is a constant depending on the mass $M$ and Newton constant $G$ that you should determine.
[8 marks]
(iii) Show that for a circular orbit, with constant radius $R=R_{0}$, then,

$$
V^{\prime \prime}\left(R_{0}\right)=\frac{J^{2}}{R_{0}^{4}}\left(1+\frac{6 \alpha}{R_{0}}\right) .
$$

[8 marks]
(iv) Compute the proper time $\tau_{\text {ang }}$ required for $\Phi$ to traverse an angle $2 \pi$. Show that for a circular orbit radius $R=R_{0}$ that is perturbed a little, so $R(\tau) \simeq R_{0}+\delta R(\tau)$, the motion approximately performs simple harmonic oscillation with period,

$$
\tau_{\mathrm{rad}}=\frac{2 \pi}{\sqrt{V^{\prime \prime}\left(R_{0}\right)}}
$$

Comment on the relation between $\tau_{\text {ang }}$ and $\tau_{\text {rad }}$.
4. (i) Consider a particle following a timelike curve $x^{\mu}(\tau)$ in a general spacetime, where $\tau$ is the particle's proper time and $v^{\mu}=d x^{\mu} / d \tau$ is its 4 -velocity. By explicit calculation show that,

$$
v^{\mu} \nabla_{\mu} v^{\alpha}=\frac{d^{2} x^{\alpha}}{d \tau^{2}}+\Gamma^{\alpha}{ }_{\sigma \rho} \frac{d x^{\sigma}}{d \tau} \frac{d x^{\rho}}{d \tau} .
$$

How is the 4 -acceleration $a^{\mu}$ related to this expression?
[8 marks]
(ii) By carefully varying the action $L_{\text {free }}$ for a free particle,

$$
L_{\text {free }}=\int d \tau \mathcal{L}_{\text {free }}, \quad \mathcal{L}_{\text {free }}=g_{\mu \nu} \frac{d x^{\mu}}{d \tau} \frac{d x^{v}}{d \tau}
$$

explicitly show that,

$$
\frac{d}{d \tau}\left(\frac{\partial \mathcal{L}_{\text {free }}}{\partial \frac{d x^{\alpha}}{d \tau}}\right)-\frac{\partial \mathcal{L}_{\text {free }}}{\partial x^{\alpha}}=2 g_{\alpha \beta}\left(\frac{d^{2} x^{\beta}}{d \tau^{2}}+\Gamma_{\sigma \rho}^{\beta} \frac{d x^{\sigma}}{d \tau} \frac{d x^{\rho}}{d \tau}\right) .
$$

[8 marks]
(iii) Suppose a particle's equation of motion is given by the variation of a different action, $L$, where,

$$
L=\int d \tau\left(\mathcal{L}_{\text {free }}+\mathcal{L}_{\text {interaction }}\right)
$$

with $\mathcal{L}_{\text {free }}$ as above and $\mathcal{L}_{\text {interaction }}$ determines the particle's interaction with some other field, so that the particle is accelerated and does not follow a geodesic. Use your previous results to deduce that the Euler-Lagrange equations of this action give rise to a particle motion with 4-acceleration given by,

$$
a_{\alpha}=-\frac{1}{2}\left(\frac{d}{d \tau}\left(\frac{\partial \mathcal{L}_{\text {interaction }}}{\partial \frac{d \alpha^{\alpha}}{d \tau}}\right)-\frac{\partial \mathcal{L}_{\text {interaction }}}{\partial x^{\alpha}}\right)
$$

[6 marks]
(iv) Consider now a particle coupled to a vector field $A_{\mu}(x)$ in a general spacetime so that its Lagrangian is modified to,

$$
L=\int d \tau\left(g_{\mu \nu}(x) \frac{d x^{\mu}}{d \tau} \frac{d x^{\nu}}{d \tau}+A_{\mu}(x) \frac{d x^{\mu}}{d \tau}\right) .
$$

Show that the 4-acceleration of the particle is;

$$
a^{\mu}=\frac{1}{2} F^{\mu v} v_{v}, \quad F_{\mu \nu}=\nabla_{\mu} A_{v}-\nabla_{v} A_{\mu}
$$

5. (i) Show that the Christoffel symbol is related to partial derivatives of the metric as,

$$
\partial_{\alpha} g_{\mu \nu}=g_{\mu \beta} \Gamma^{\beta}{ }_{\alpha \nu}+g_{\nu \beta} \Gamma^{\beta}{ }_{\alpha \mu} .
$$

(ii) The Lie derivative of a $(0,2)$ tensor $A_{\mu \nu}$ with respect to a vector field $w^{\mu}$ is,

$$
(\text { Lie })(w, A)_{\mu \nu}=w^{\alpha} \partial_{\alpha} A_{\mu \nu}+A_{\mu \alpha} \partial_{\nu} w^{\alpha}+A_{\alpha \nu} \partial_{\mu} w^{\alpha} .
$$

Suppose we consider the Lie derivative of the metric $g_{\mu v}$. Show that this can also be written in terms of the covariant derivative as,

$$
(\text { Lie })(w, g)_{\mu \nu}=\nabla_{\mu} w_{v}+\nabla_{\nu} w_{\mu} .
$$

If this vanishes, we say $w^{\mu}$ is a Killing vector field.
[6 marks]
(iii) Consider a timelike particle with velocity $v^{\mu}=d x^{\mu} / d \tau$ for proper time $\tau$. Suppose it follows a geodesic in a spacetime with a Killing vector field $w^{\mu}$. Show that the quantity,

$$
\phi=-w^{\mu} v_{\mu}
$$

is constant along the particle's trajectory.
(iv) Consider the spacetime with coordinates $x^{\mu}=\left(t, x^{i}\right)$

$$
d s^{2}=-N(x) d t^{2}+g_{i j}(x) d x^{i} d x^{j}
$$

where $N$ and $g_{i j}$ only depend on the spatial coordinates $x^{i}$ and not time $t$. Show that there is a Killing vector $w^{\mu}$ for this spacetime and explicitly check that $\operatorname{Lie}(w, g)=0$. In this spacetime how is the conserved quantity $\phi$ for a non-accelerated particle related to the energy of the particle as measured by observers sitting at constant spatial position?
[10 marks]
[Total 30 marks]

