

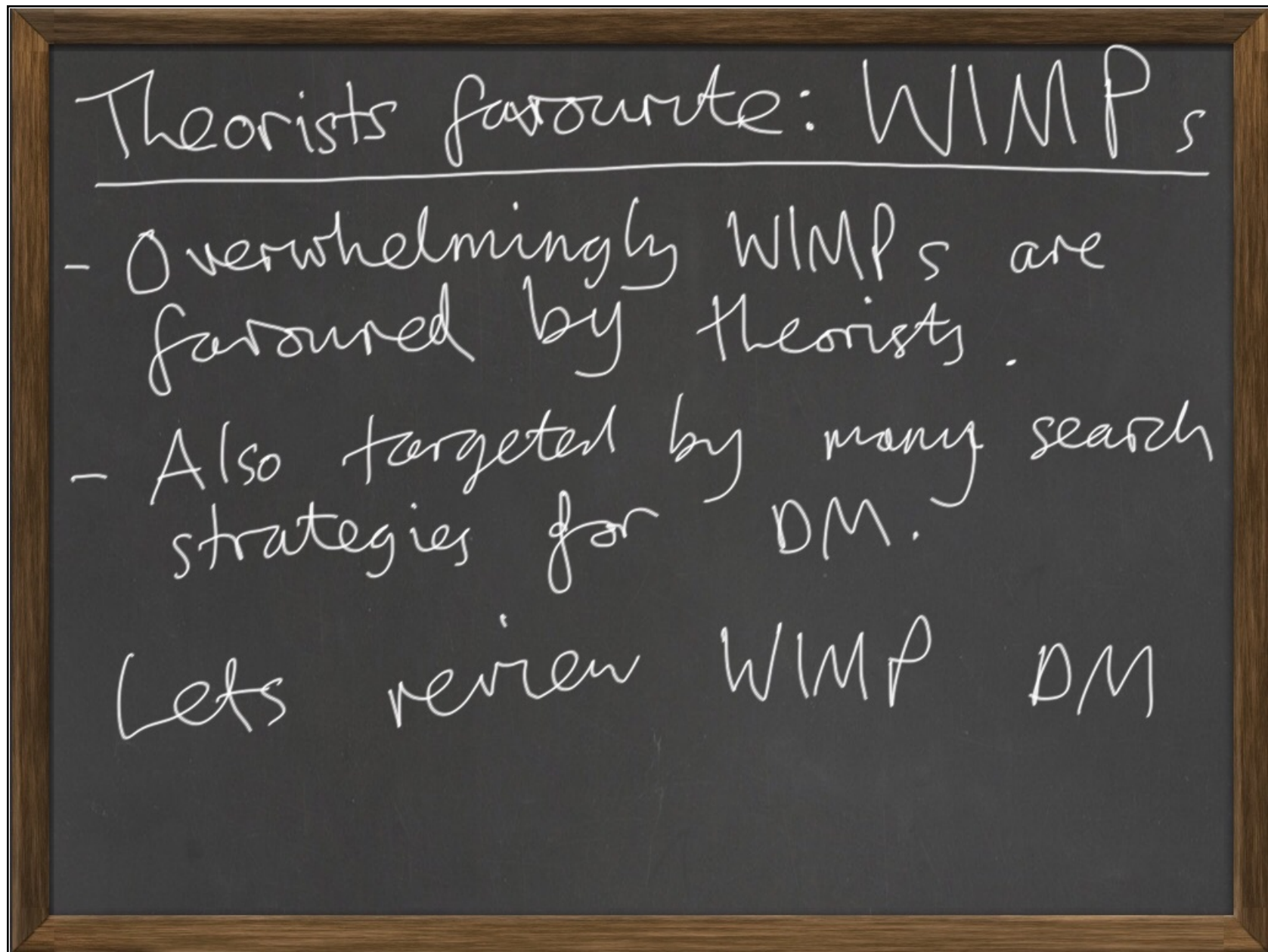
Based on work done with

• G. Kane, P. Kumar, K. Bobkov, S. Watson.
(Non-thermal)

• S. Ellis, G. Kane, B. Nelson, M. Perry
(DM is Hidden)

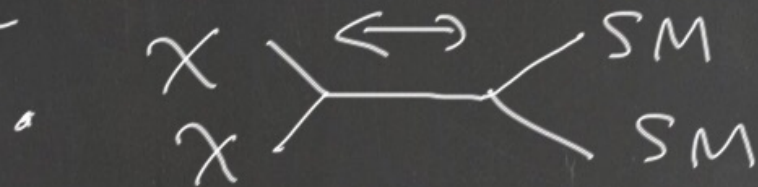
arXiv 1604.05320, PRL 117, 181802, 2016.

• M. Fairbairn, E. Hardy, arXiv 1704.01804



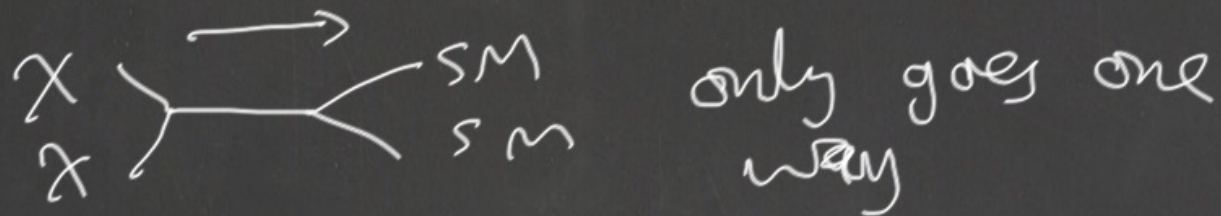
At the end of inflation (or whatever solves the horizon, flatness probs and seeds the CMB!):

- Assume Universe is radiation dominated with a high $T \gg M_{EW} \sim 100 \text{ GeV}$
- Standard Model particles are in equilibrium with WIMPs, X



X is a stable, electrically neutral particle charged under $SU(2) \times U(1)_Y$.

As Universe expands, T drops.
 When T falls below m_X ,



And X particles freeze out with

$$H \Big|_{T \sim \frac{m_X}{\text{few}}} \sim n_X$$

$$\langle \sigma v \rangle_{XX \rightarrow SM} < (6v)$$

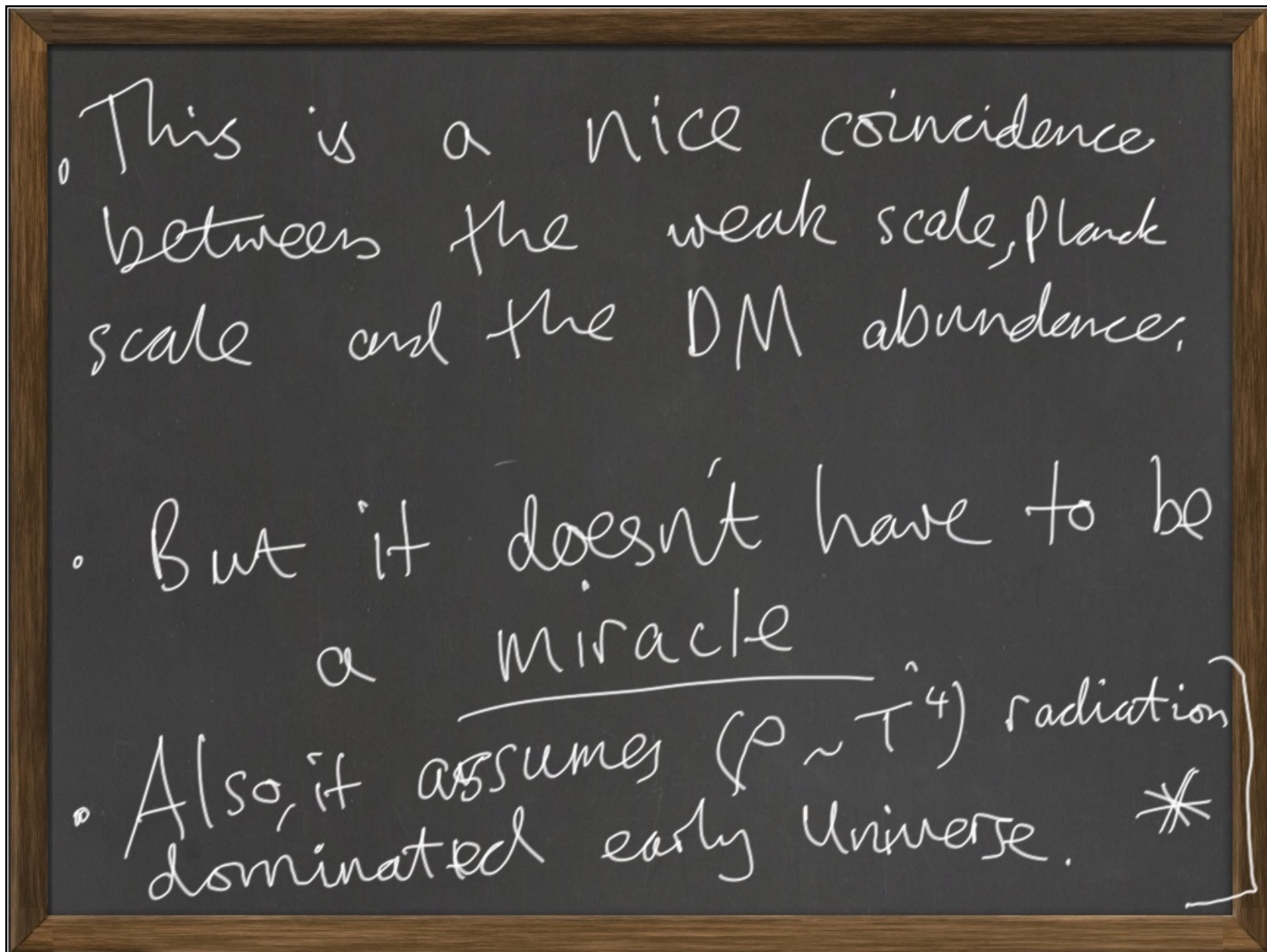
$$H \sim \frac{T^2}{m_{pl}} \quad G_U \approx \frac{\alpha^2}{M_x^2}$$

so $\nu_x \approx \frac{T^2 M_x^2}{\alpha^2 m_{pl}}$

$\frac{\rho}{S} \approx \frac{M_x^3}{g_* \alpha^2 T m_{pl}}$ $\left(S \sim g_* T^3 \right)$

If $T \sim \frac{M_x}{S}$
 $\alpha \sim 10^{-2} S$

$\frac{\rho}{S} \sim 10^5 \frac{M_x^2}{m_{pl}} \sim$
 $\sim 10^{-10} \text{ GeV} \checkmark$



We will consider the low energy limits of solutions of string/M-theory

\exists many solutions of the form:

$$M^{9,1} = \underbrace{Z^6}_{\text{compact, small}} \times \underbrace{M^{3,1}}_{\text{large}} \quad \leftarrow \text{string}$$

or

$$M^{10,1} = X^7 \times M^{3,1} \quad \leftarrow \text{M theory eg } G_2.$$

$$\underline{g(M^{10,1})} \cong \boxed{g(x)} + \underline{g(M^{3,1})} \quad \leftarrow$$

Low energy, $d=3+1$ Lagrangian is of the form, schematically,

$$\begin{aligned}
 - \int_{\text{matter} + \text{grains}} &= \left. \begin{aligned} &\frac{1}{16\pi G_N} \sqrt{-g_{3+1}} R_{3+1} + \frac{1}{g^2} F_{\mu\nu}^2 \\ &+ i \bar{\Psi} \not{\partial} \Psi + \lambda H \bar{\Psi} \Psi \\ &+ |\mathcal{D}H|^2 - V(H, H^\dagger) \end{aligned} \right\} \\
 + \\
 \underline{\underline{\int_{\text{moduli}}}} &= \underline{\underline{\kappa^{ij}(s_i)}} \left(\underline{\underline{d_\mu s_i d^\mu s_j}} + \underline{\underline{\kappa^{ij}(s) d_\mu a_i d^\mu a_j}} \right) \\
 &\quad - \underline{\underline{V(s_i, a_j)}} \\
 s_i &= \text{moduli} \quad a_i = \text{axions} \quad + \dots
 \end{aligned}$$

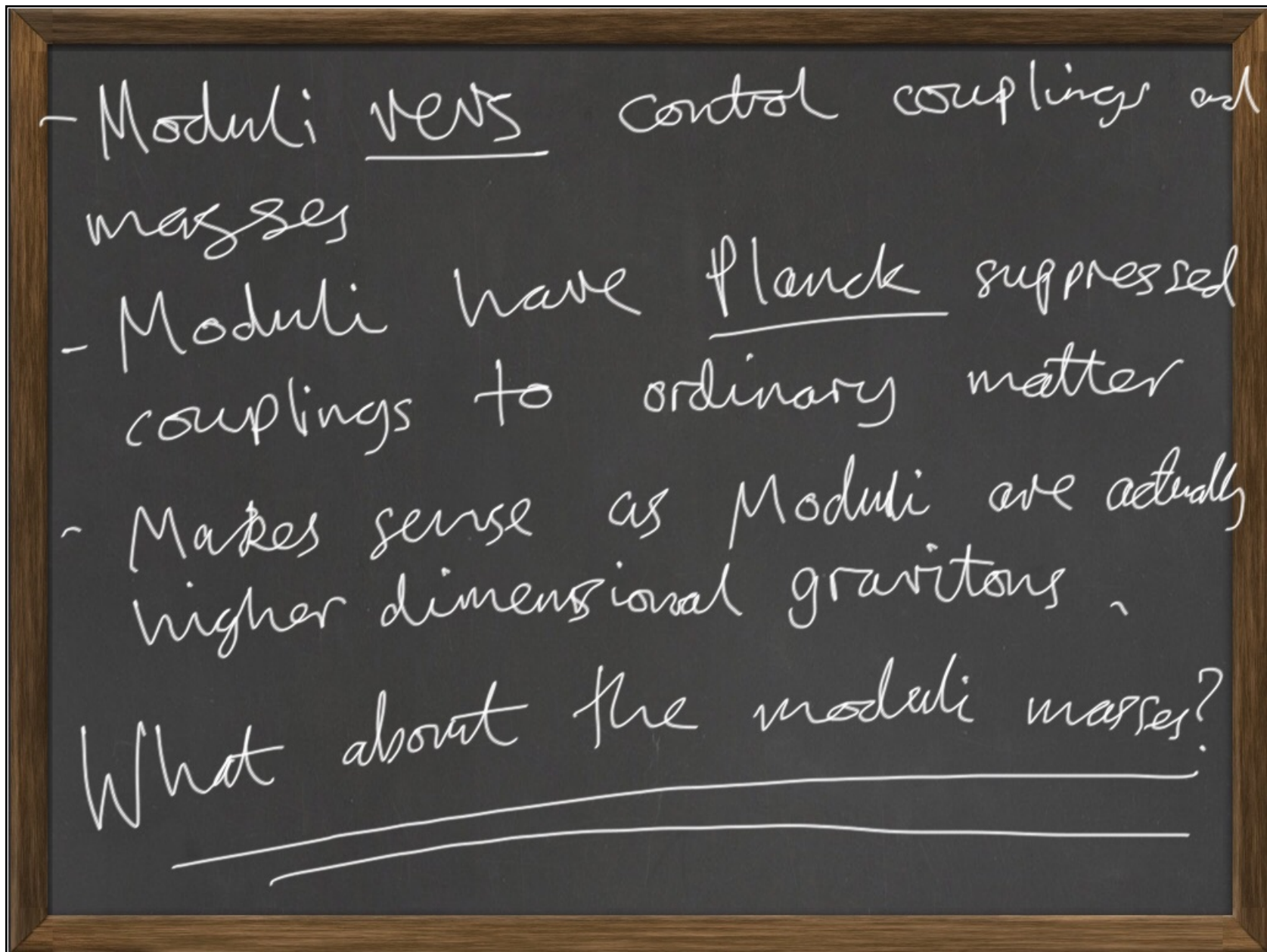
Moreover $\frac{1}{g^2} F_{\mu\nu}^2$ is really

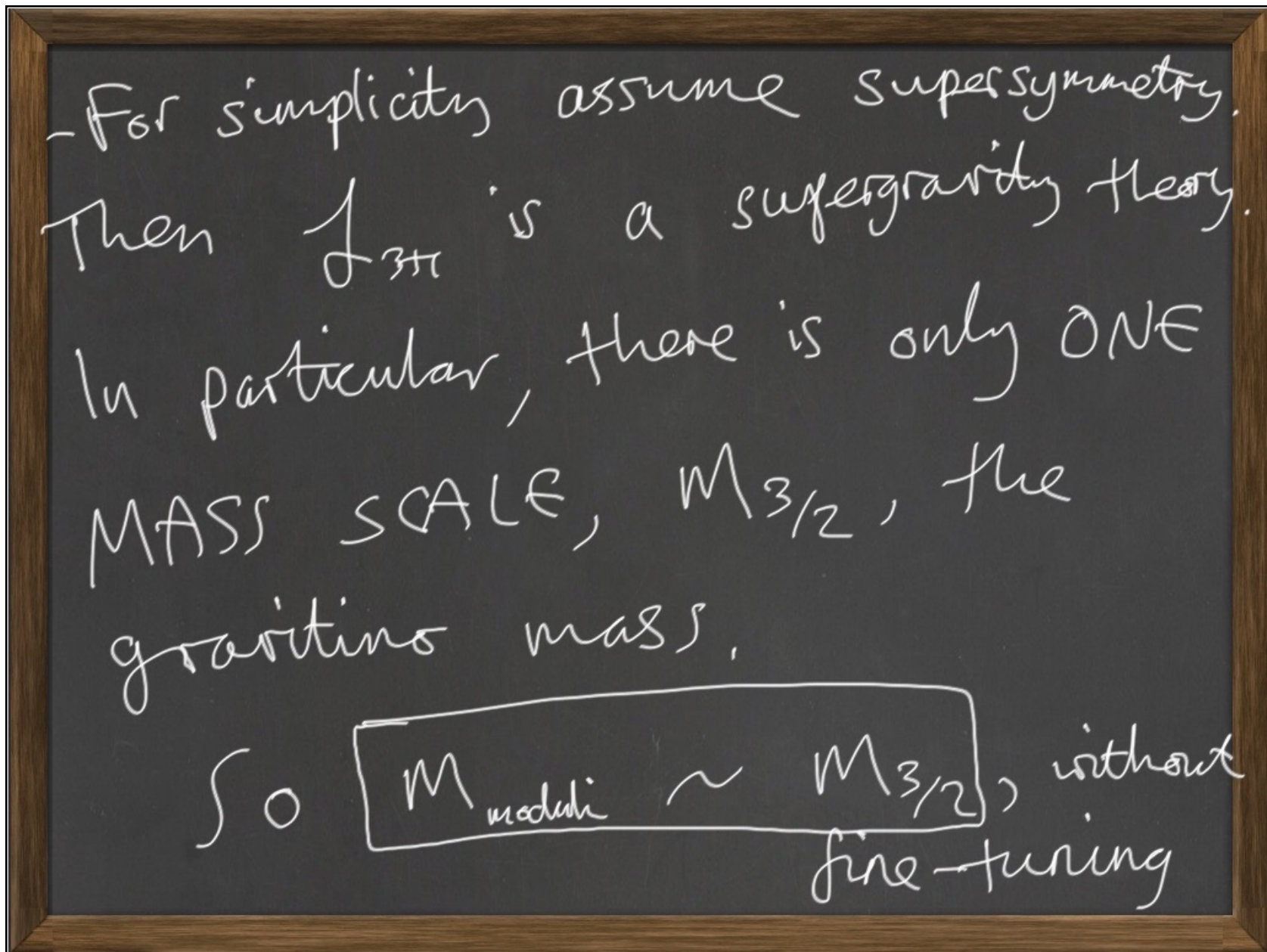
$S \sqrt{g} F_{\mu\nu}^2$ $\frac{N_i S_i F_{\mu\nu}^2}{m p^2}$, a Dirac operator.

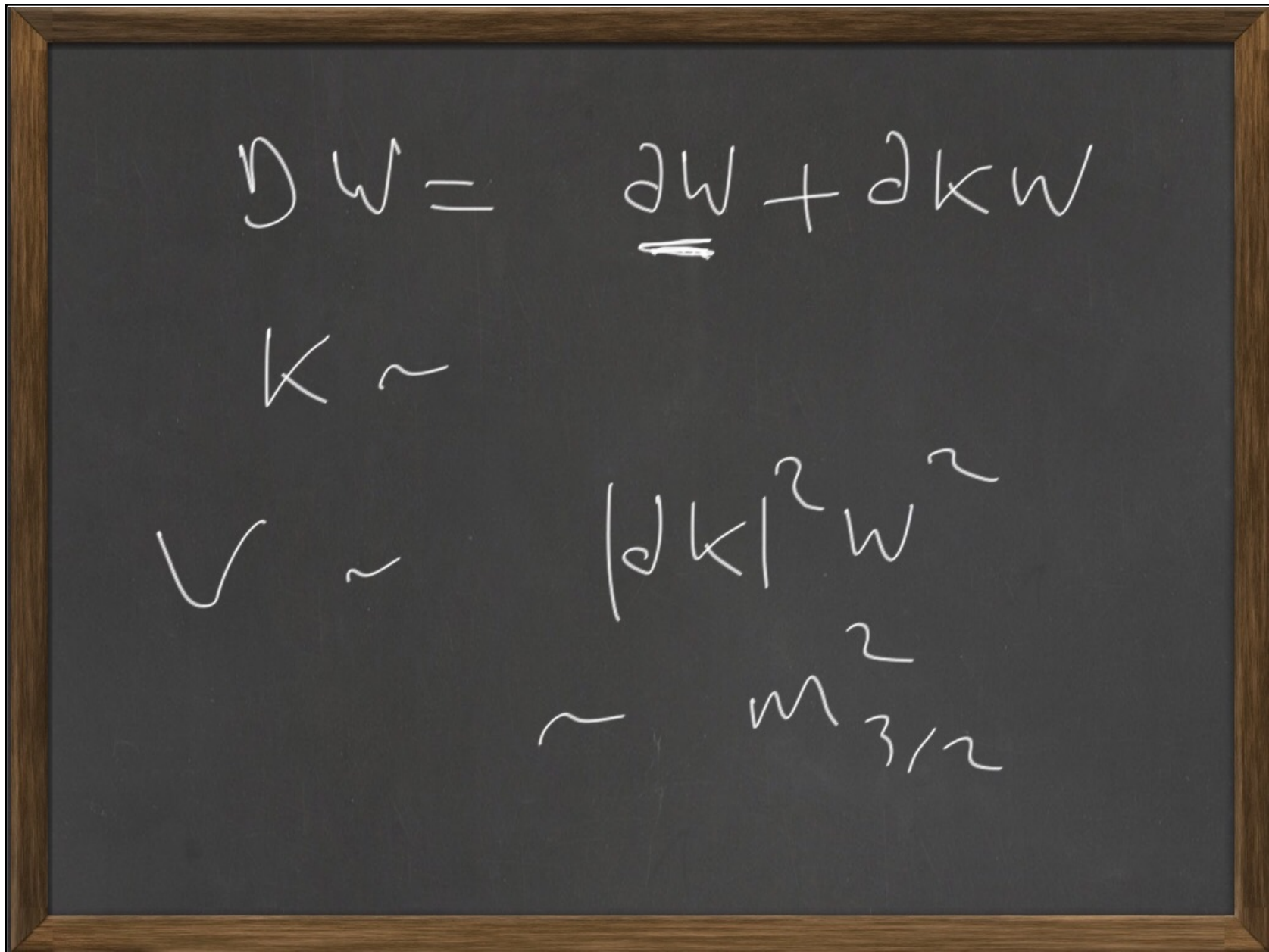
Similarly $\frac{\lambda}{\Lambda} \psi_L \psi_R$ is really

$\psi \rightarrow e^{-\frac{d_i S_i}{m p^2}} e^{\frac{i d_i a_i}{f}} \psi_L \psi_R$

* The moduli dependence of λ varies from theory to theory.







Cosmology of This Theory

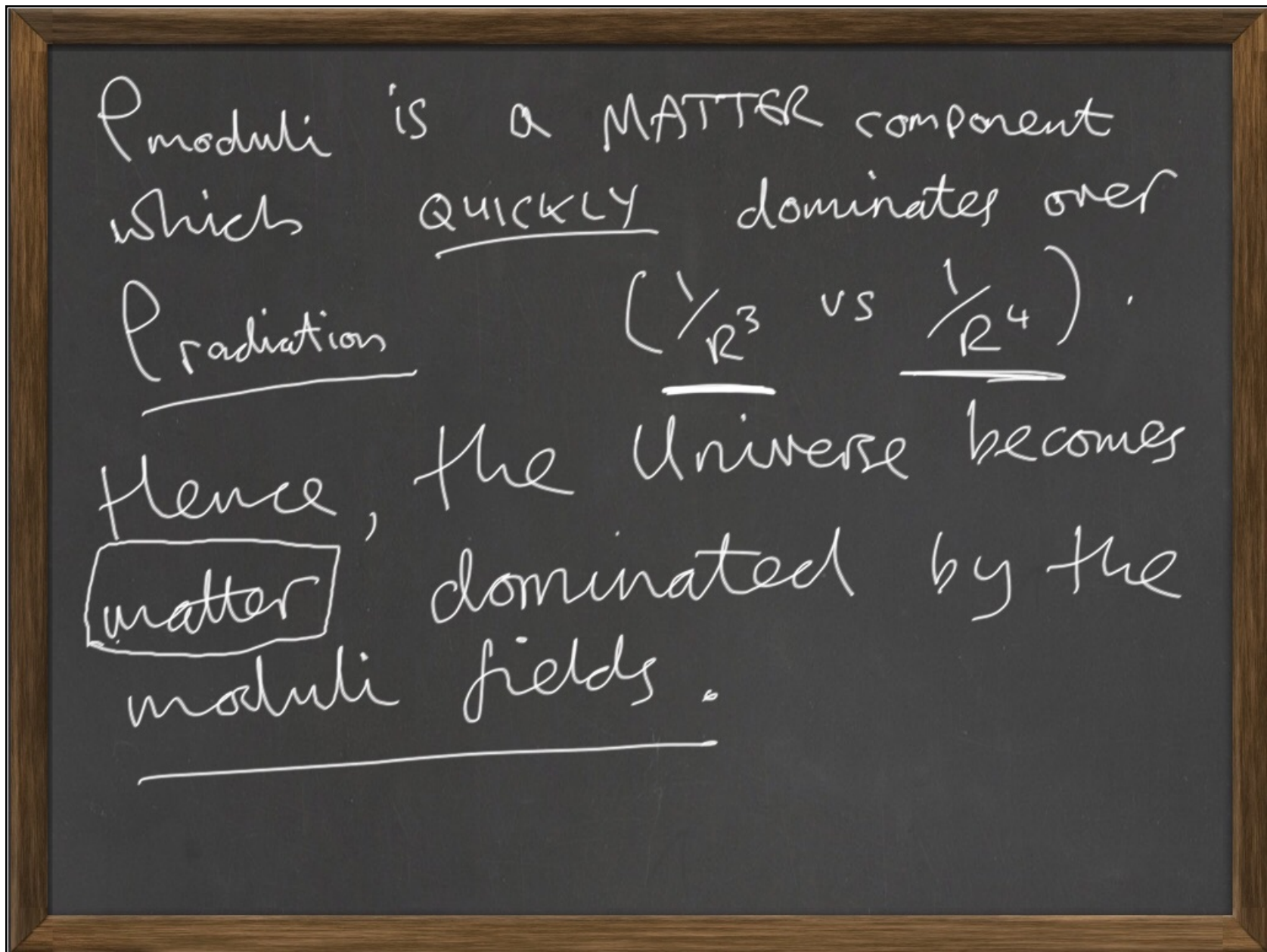
At the end of inflation (or whatever.....)

If $H \gg m_{3/2} \sim m_{\text{moduli}}$, the moduli will be stuck at some $O(1) m_{\text{pl}}$ place in its potential.

Later, $H \sim O(m_{\text{moduli}})$ and s_i oscillates:

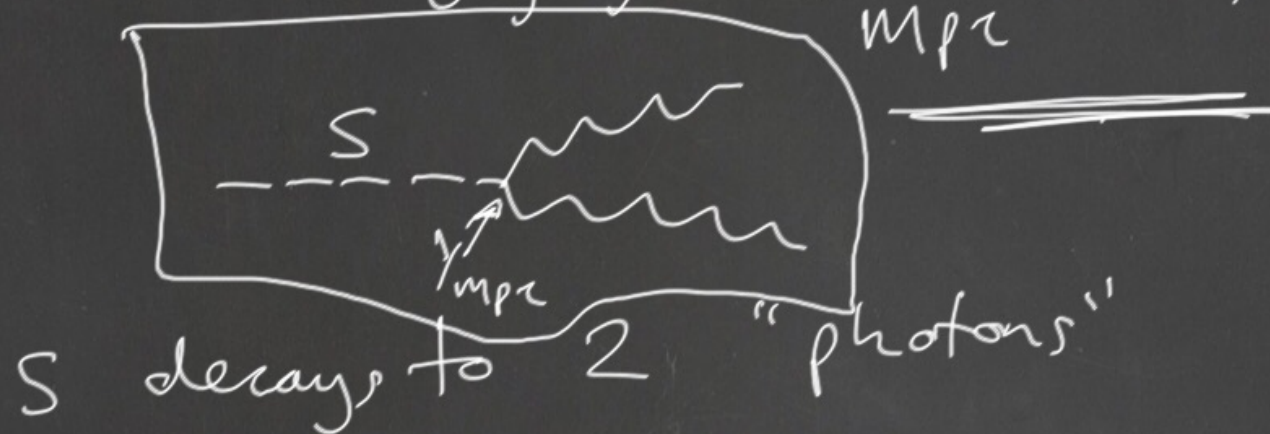
The top diagram is a plot of potential energy $V(s)$ versus s . It shows a curve with a local minimum. A point on the curve is marked with a dot and labeled m_{pl} , with a horizontal double-headed arrow indicating the width of the potential well at that level.

The bottom diagram is a plot of potential energy $V(\text{mod})$ versus mod . It shows a curve with a local minimum. A point on the curve is marked with a dot and labeled $H \sim O(m_{\text{mod}})$, with a horizontal double-headed arrow indicating the width of the potential well at that level.



The moduli are unstable particles.
 (They couple to matter particles fairly
 'generically' and 'uniformly'.)

(Consider, e.g., $\mathcal{L}_{\text{gauge}} \sim \frac{S}{M_{\text{pl}}^2} F_{\mu\nu}^2$



Decay width (or probability) is

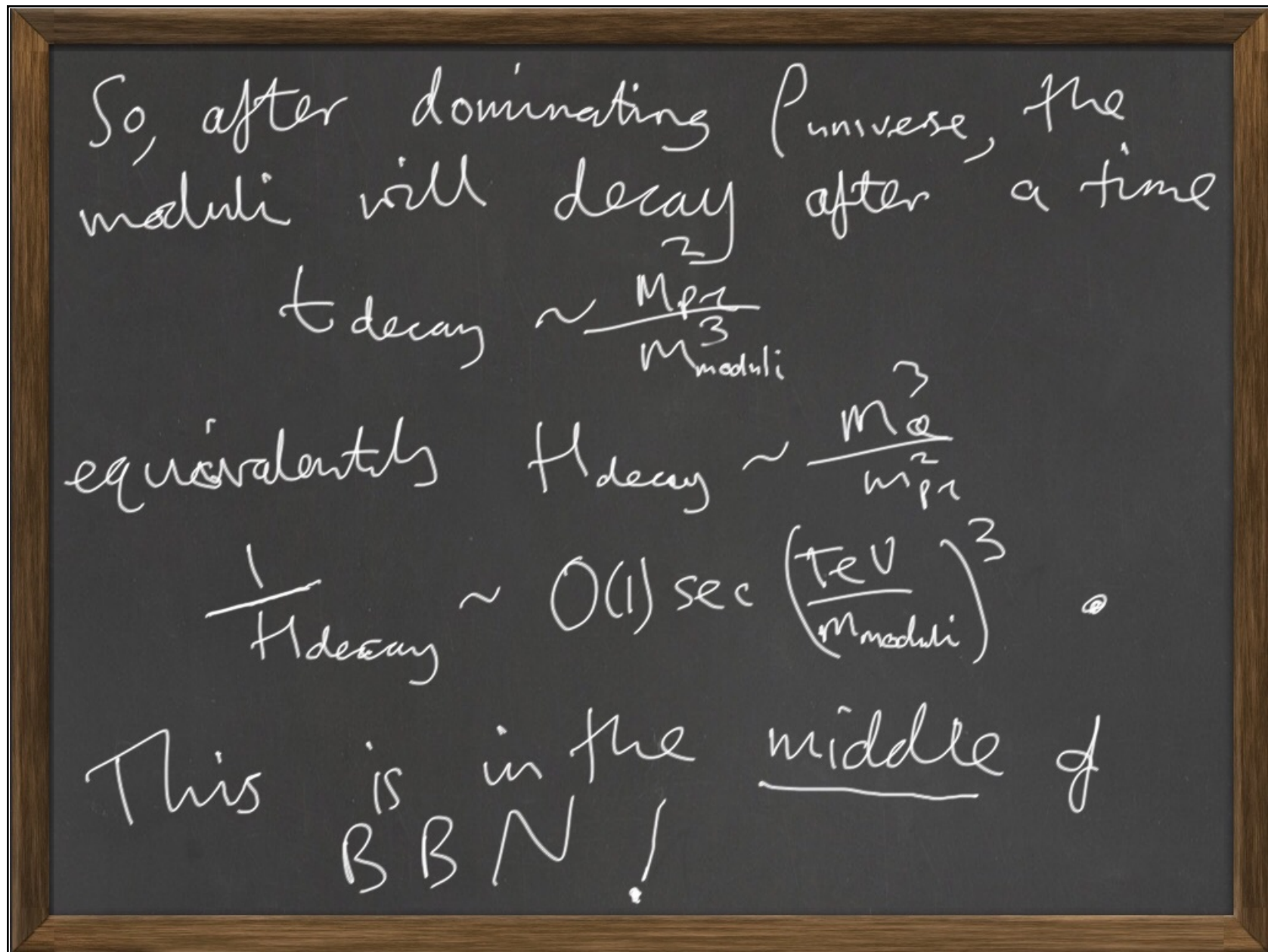
$$\Gamma(S \rightarrow \gamma\gamma) \propto |M|^2$$

$$M \sim \frac{1}{m_{Pl}^2}$$

$$\Gamma \sim O\left(\frac{1}{m_{Pl}^2}\right) \sim G_N \text{ (for Easy!)}$$

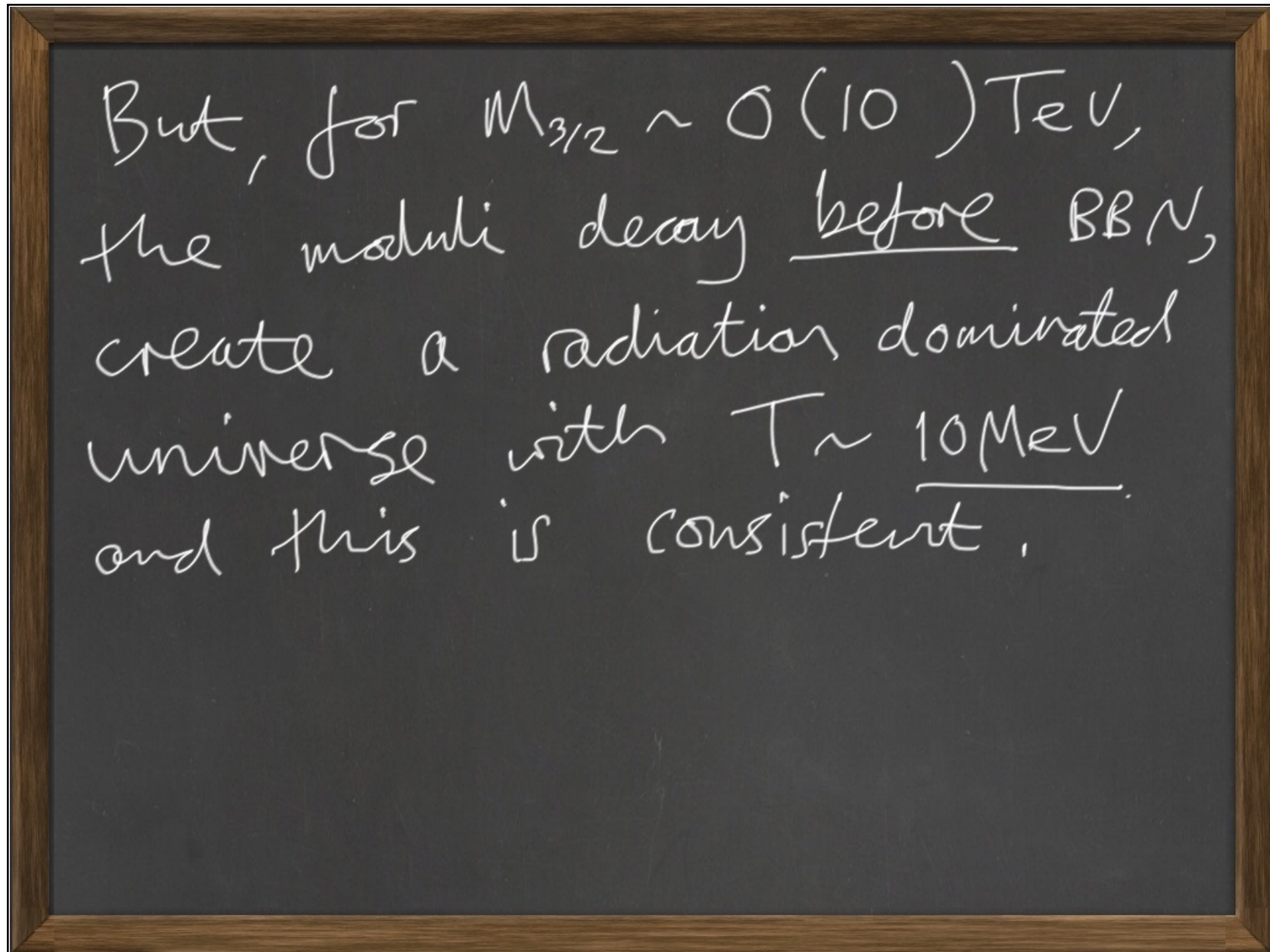
$$\therefore \Gamma(S \rightarrow \gamma\gamma) \approx \frac{M_{moduli}^3}{2 m_{Pl}^2}$$

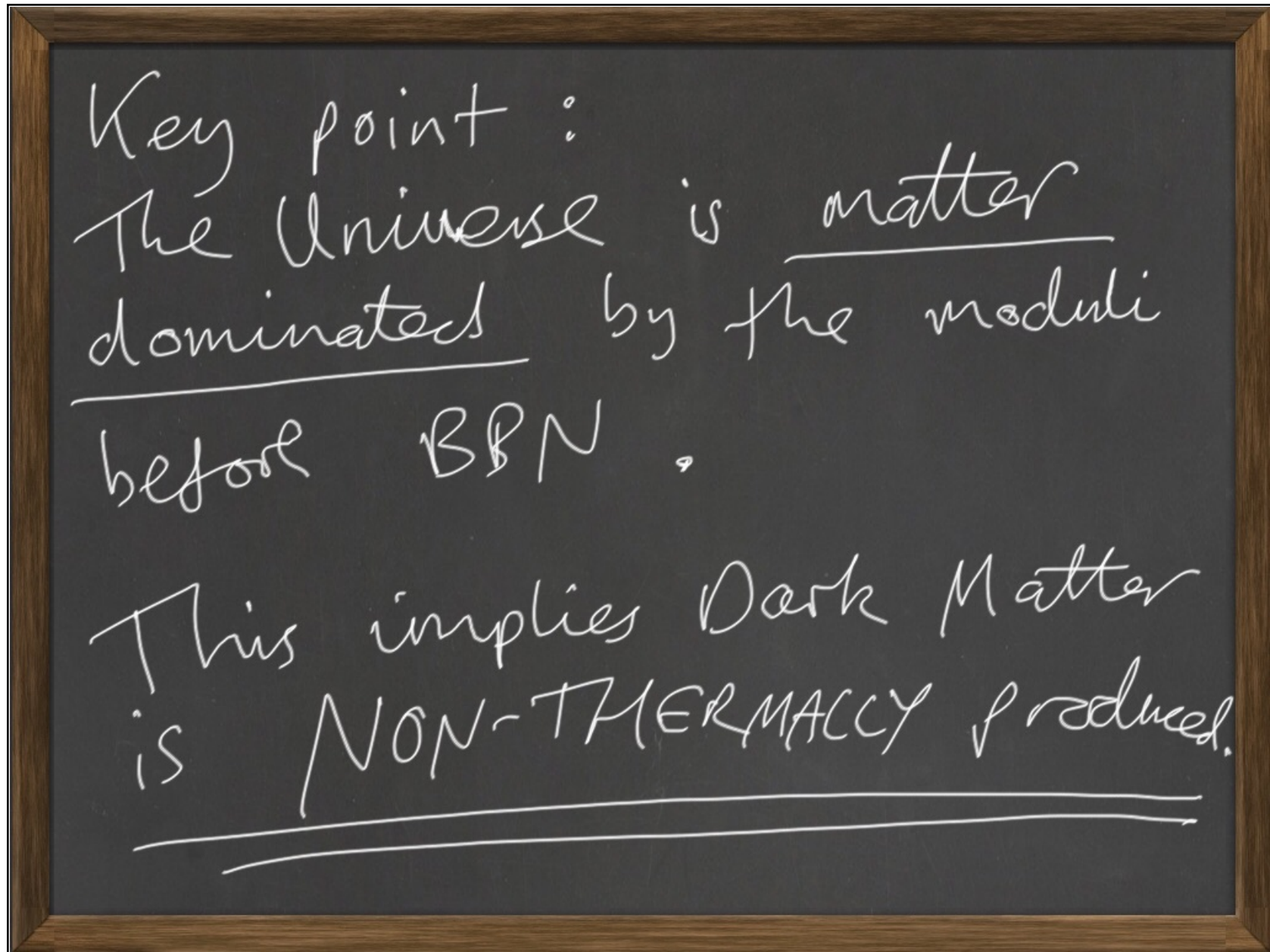
Lifetime $\tau_{moduli} \approx \frac{m_{Pl}^2}{M_{moduli}^3} \sim$



So for $M_{3,2} \sim \text{TeV}$, moduli decay during BBN. This is bad as they decay into quarks, leptons and gauge bosons.

This injects charged particles and hadrons into the plasma which can dis-associate nuclei and drastically change the successful predictions of BBN.

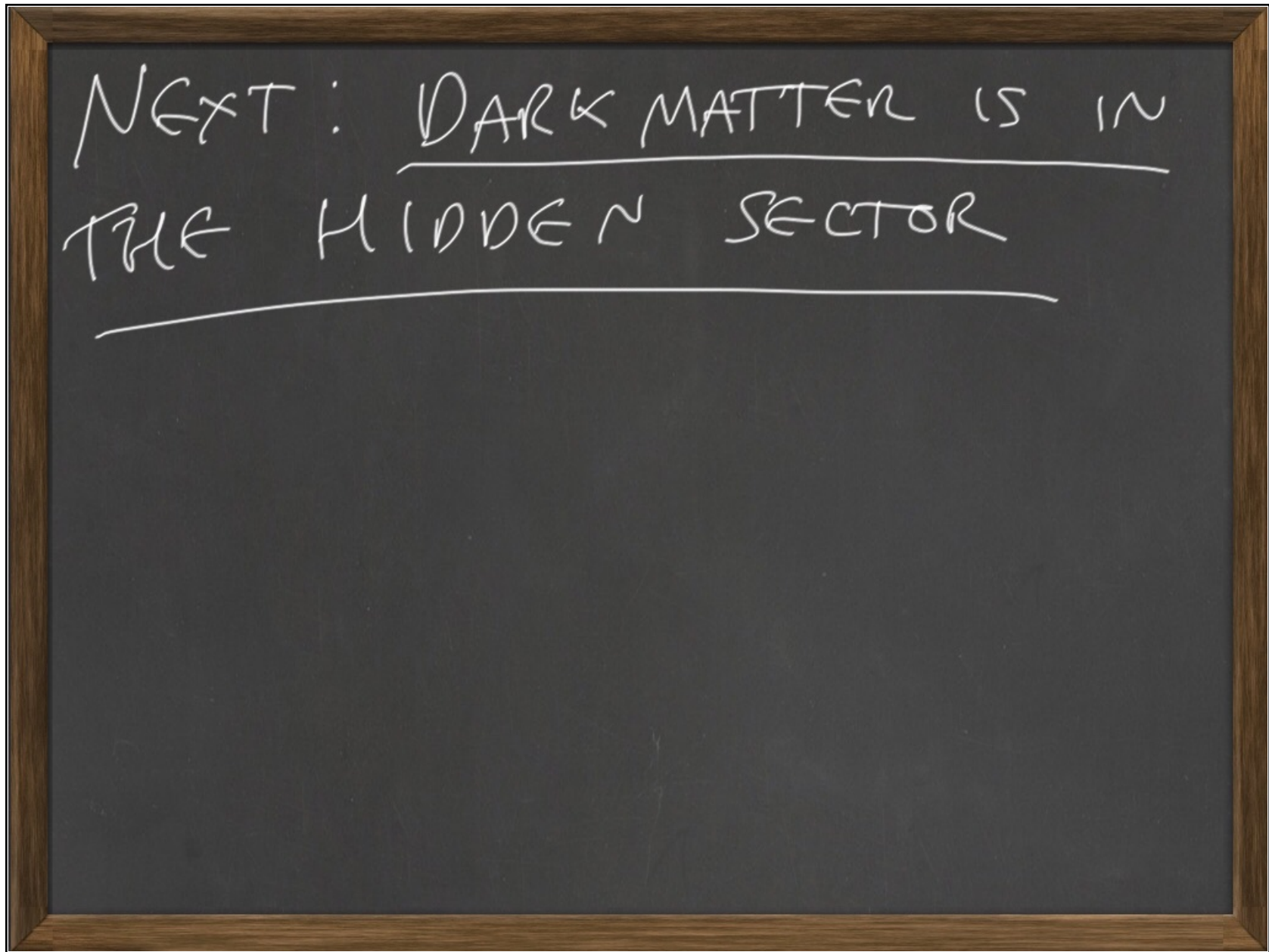




This seems quite a generic
conclusion.

Caveats

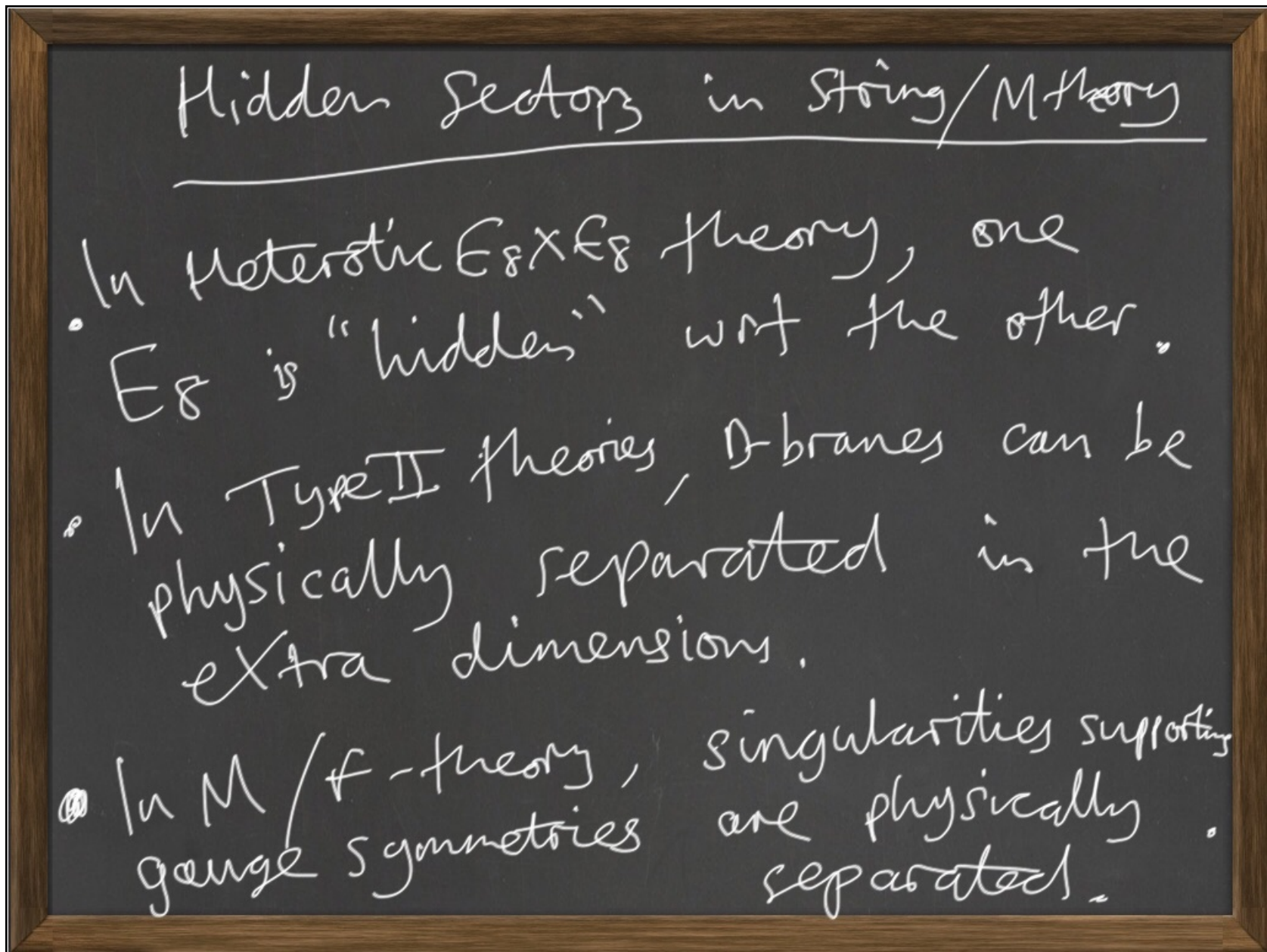
- Could assume $H_{\text{inf}} \ll M_{3/2}$
(not typical)
- Could arrange a late period
of inflation to "get rid of the
moduli". (Seems 'tuned'.)



Hidden Sectors

Defⁿ: A particle is in the
 Hidden sector if it has no
tree level gauge interactions
 with the Standard Model.
 ie it has no $SU(3) \times SU(2) \times U(1)_Y$
 charge at tree level.

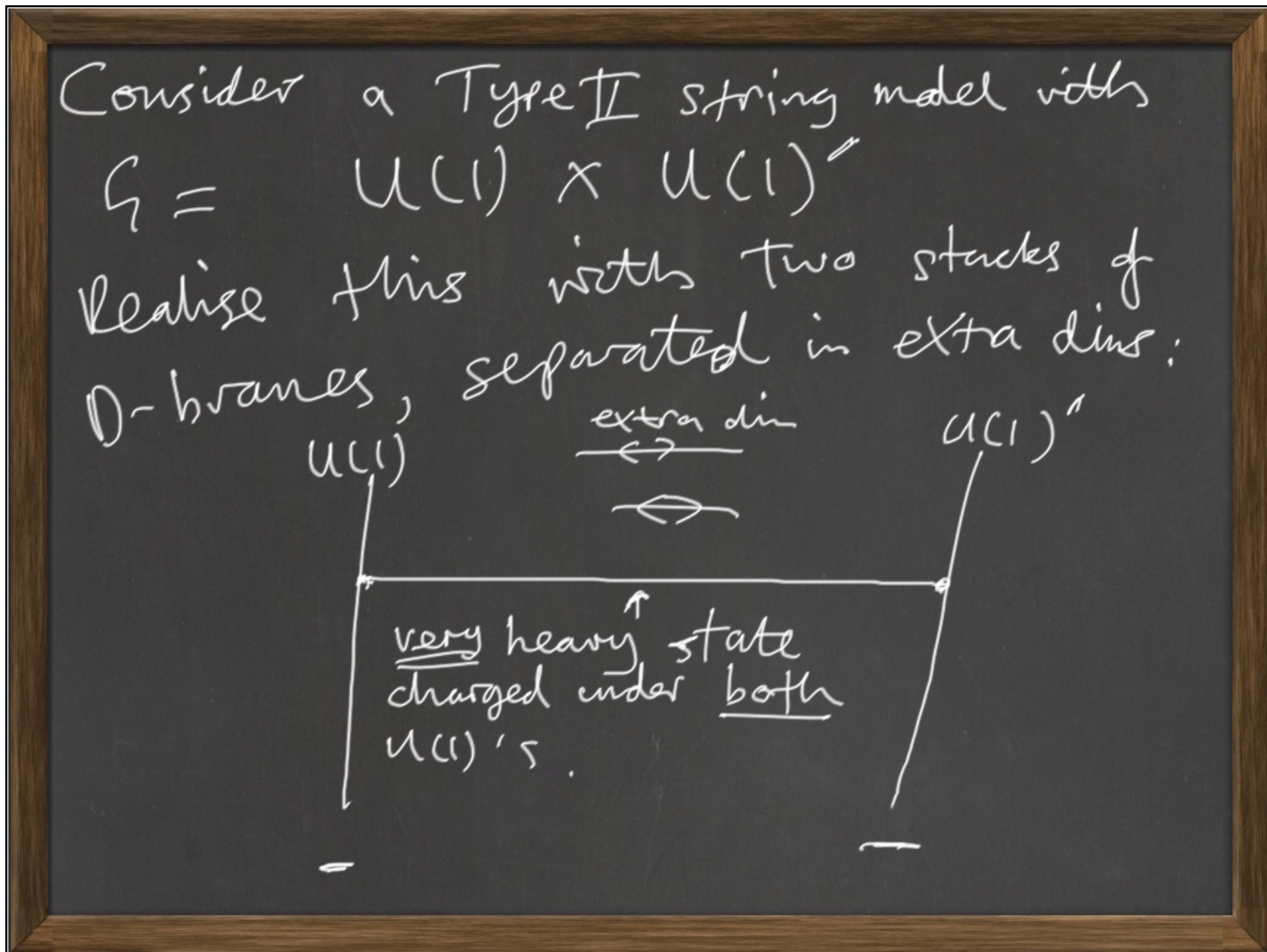
Since we have no idea why the Standard Model has $G = SU(3) \times SU(2) \times U(1)$ and 45 fermions and a Higgs doublet, there is no reason NOT to consider additional gauge sectors and matter. This is exactly the picture that emerges from string/M theory



There is no privilege given to the Standard Model.

Genericcally expect additional gauge groups and matter.

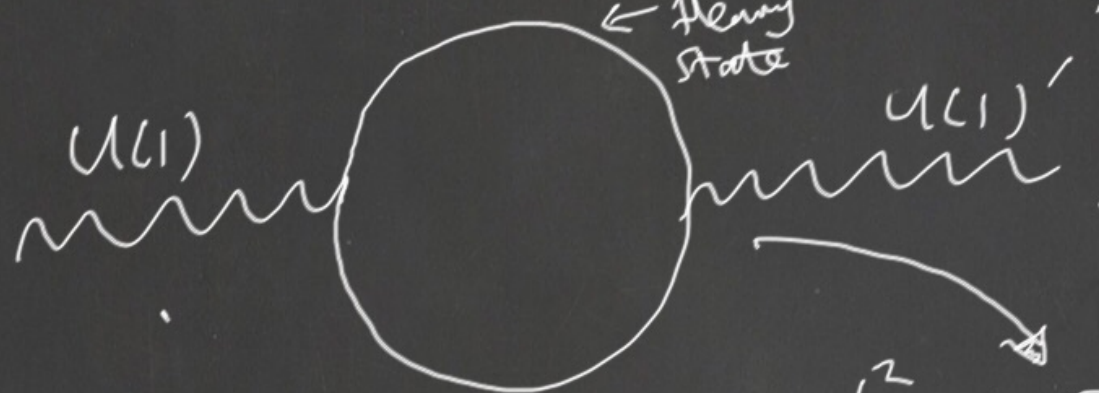
HIDDEN SECTOR MATTER
IS GENERIC



Mass, heavy state $\sim \frac{M_{\text{str}}^2 R_{\text{KK}}}{M_{\text{KK}}}$

$\gg \frac{M_{\text{KK}}}{M_{\text{KK}}}$

H induces a renormalisation of kinetic terms:



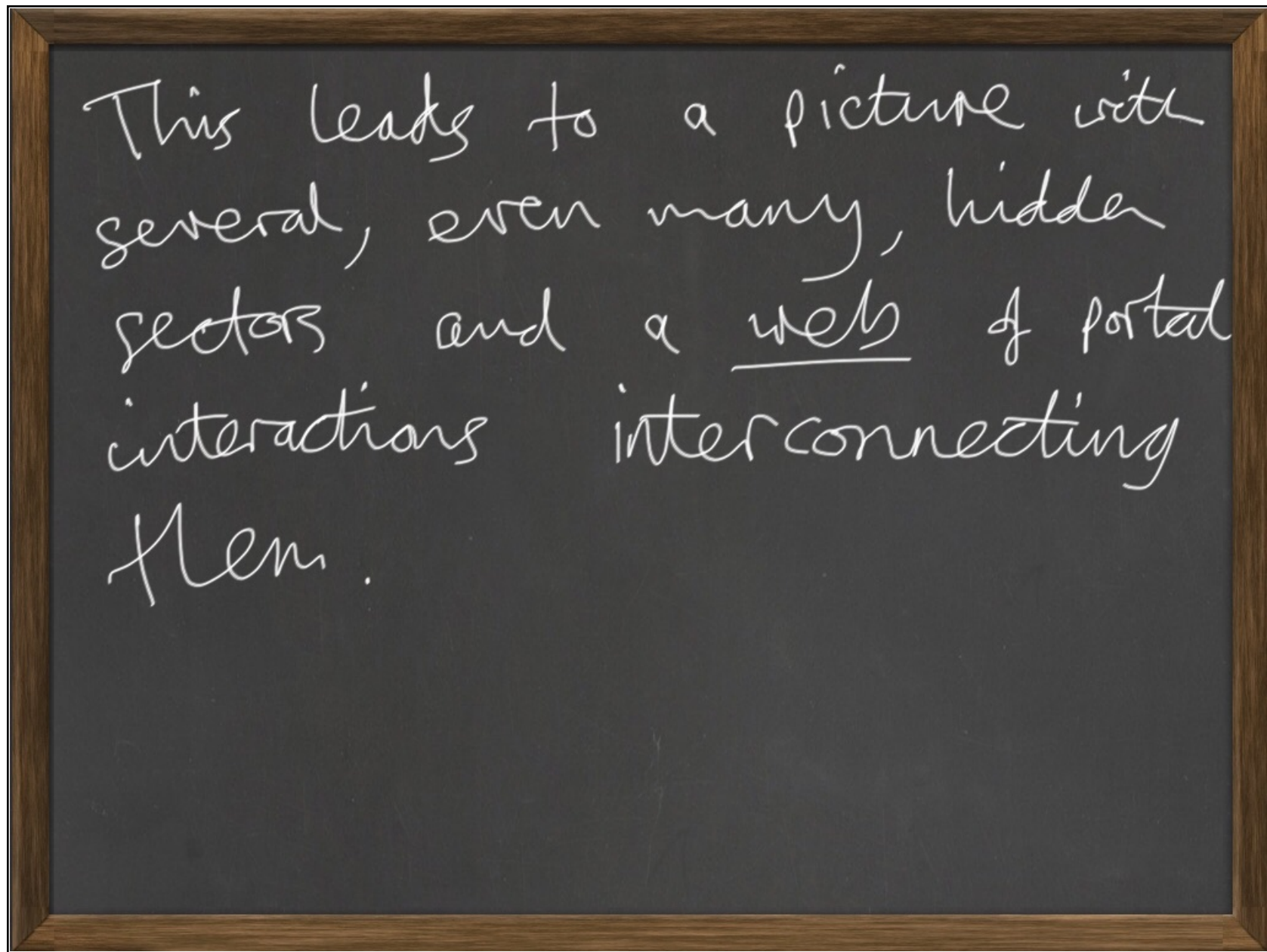
i.e. $F_{\mu\nu}^2 + \tilde{F}_{\mu\nu}^2 \rightarrow F_{\mu\nu}^2 + F_{\mu\nu}'^2 + \epsilon F_{\mu\nu} F_{\mu\nu}'$

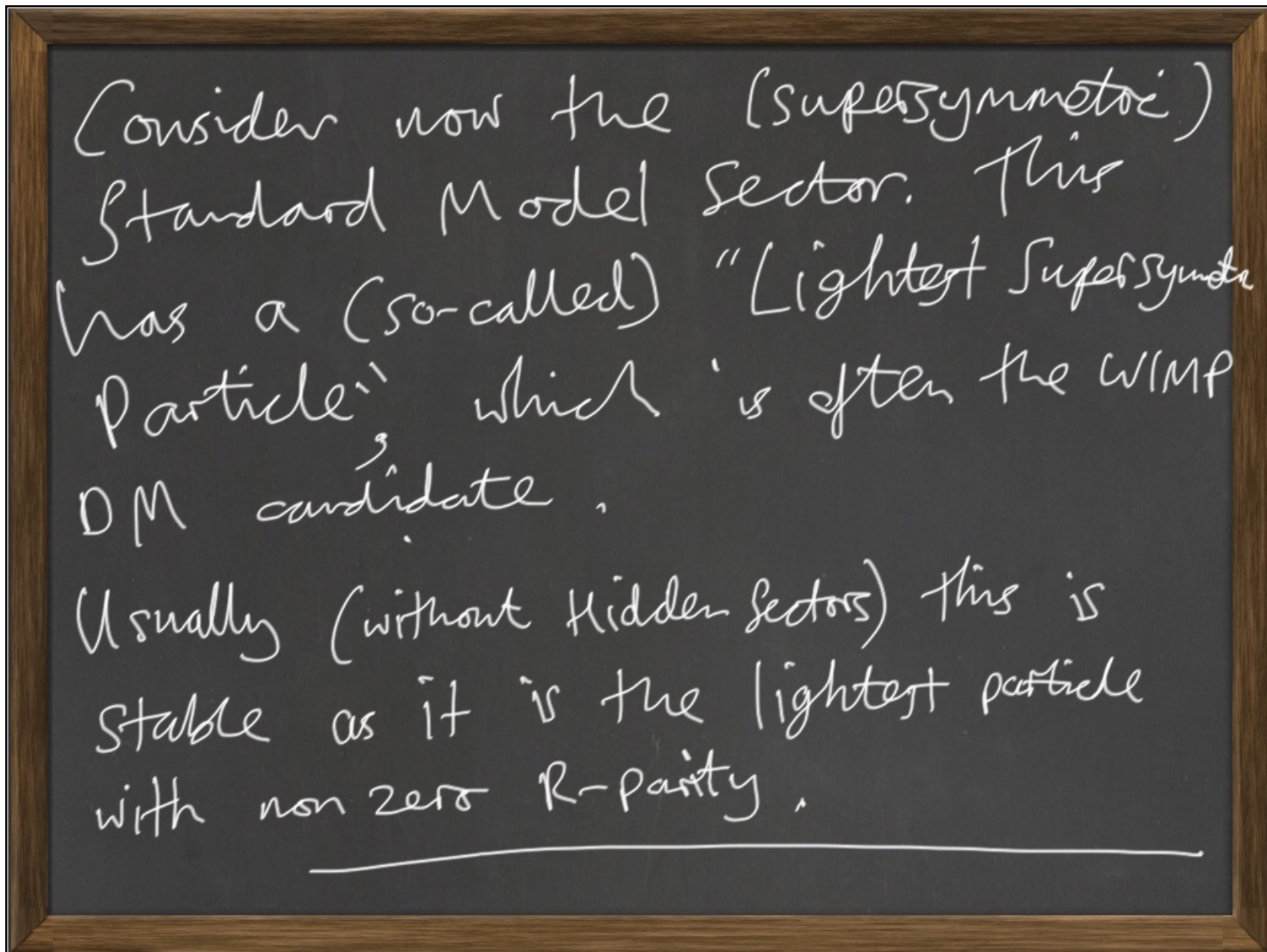
Since FF' is dim 4, ϵ is only log sensitive to UV

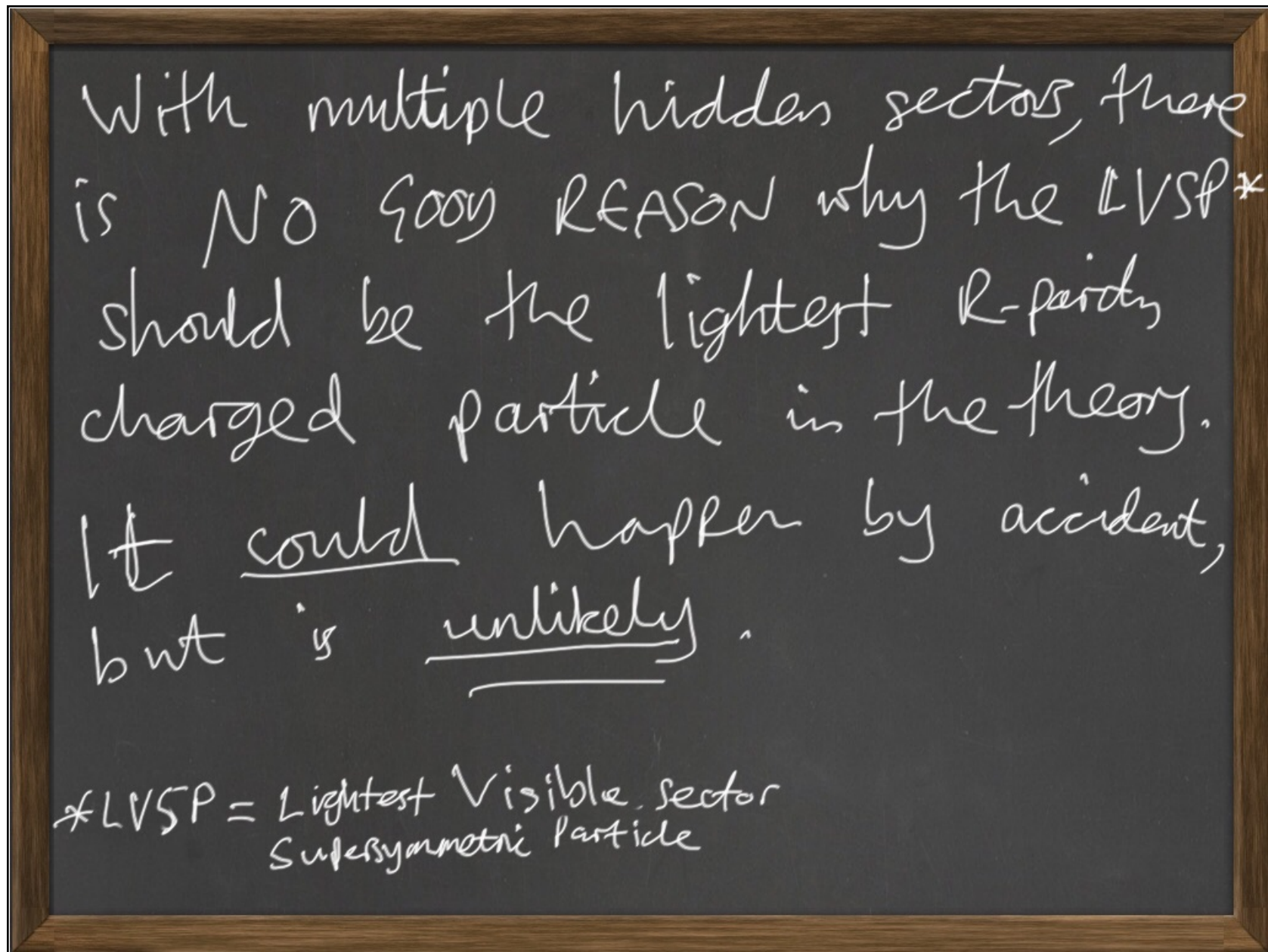
$$E \sim \frac{gg'}{12\pi^2} \ln\left(\frac{\Lambda}{M}\right).$$

- Such mixings are generically present between U(1)'s.
- This has been known for quite some time (Dienes, Kolda, March-Russell '97)

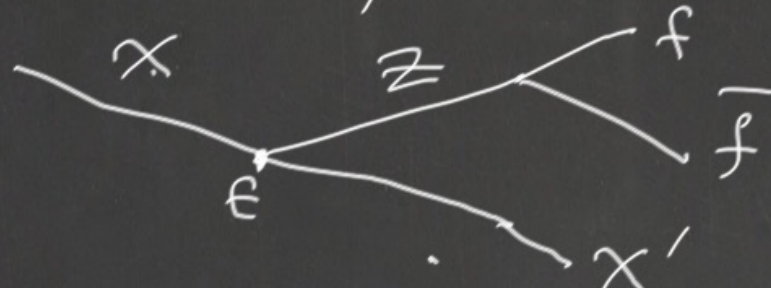
The $\epsilon F F'$ interactions (and those related to it by supersymmetry) provides a PORTAL between different hidden sectors, eg gauge bosons can mix between sectors, as can gauginos, via $\epsilon \lambda \phi \lambda'$.





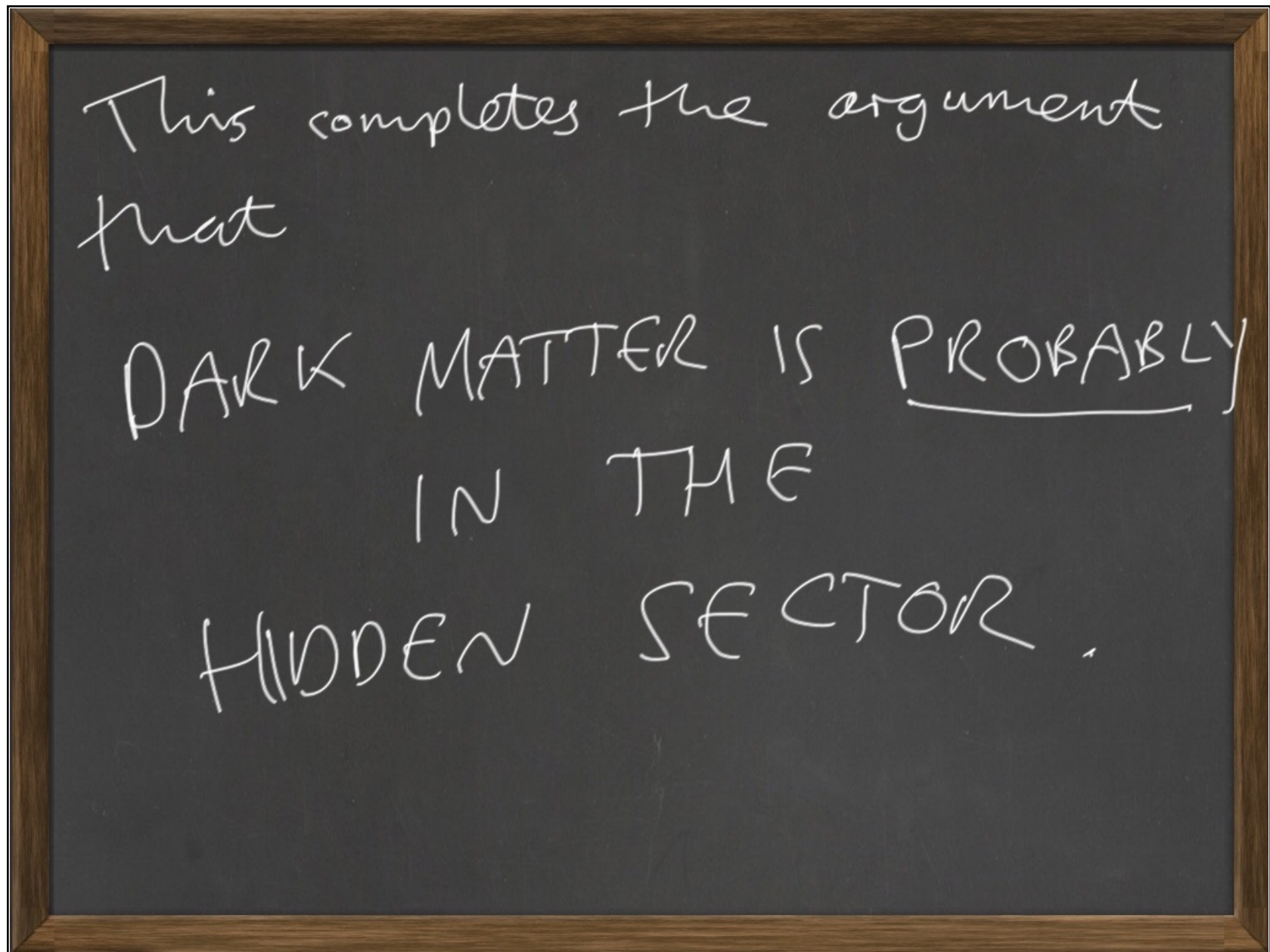


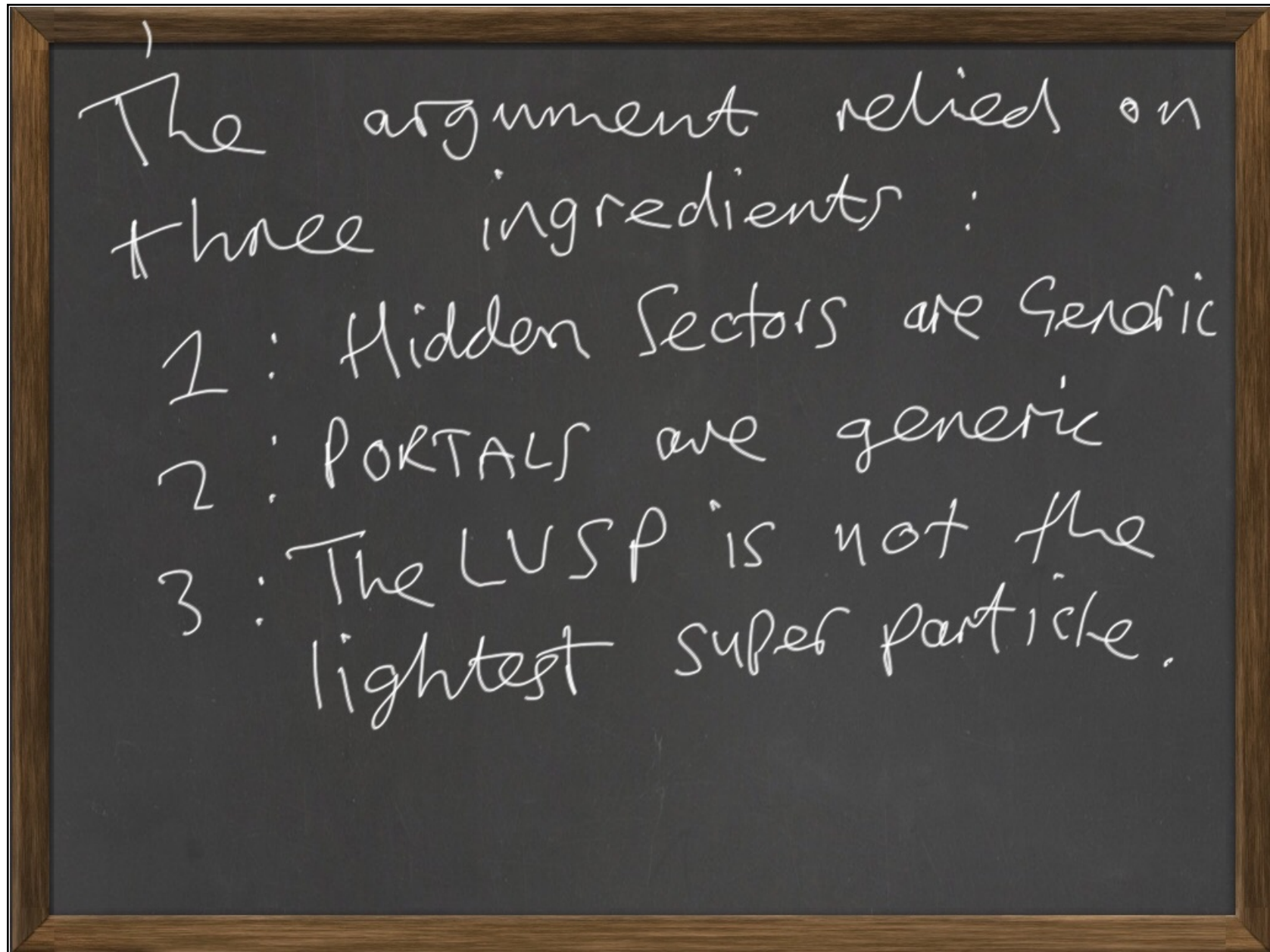
Mixing between Hidden $U(1)'$
and $U(1)_Y$ leads to, e.g.

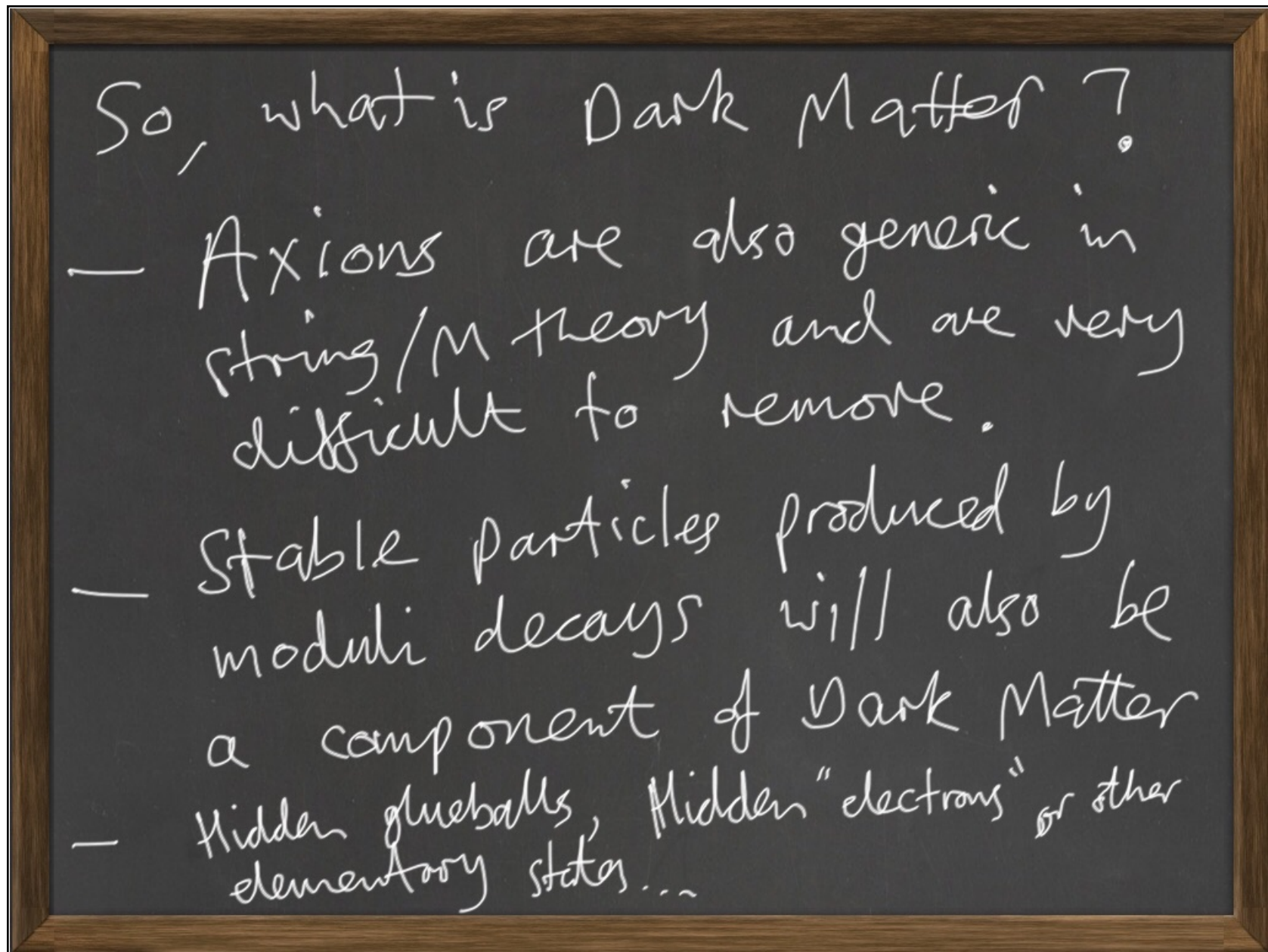


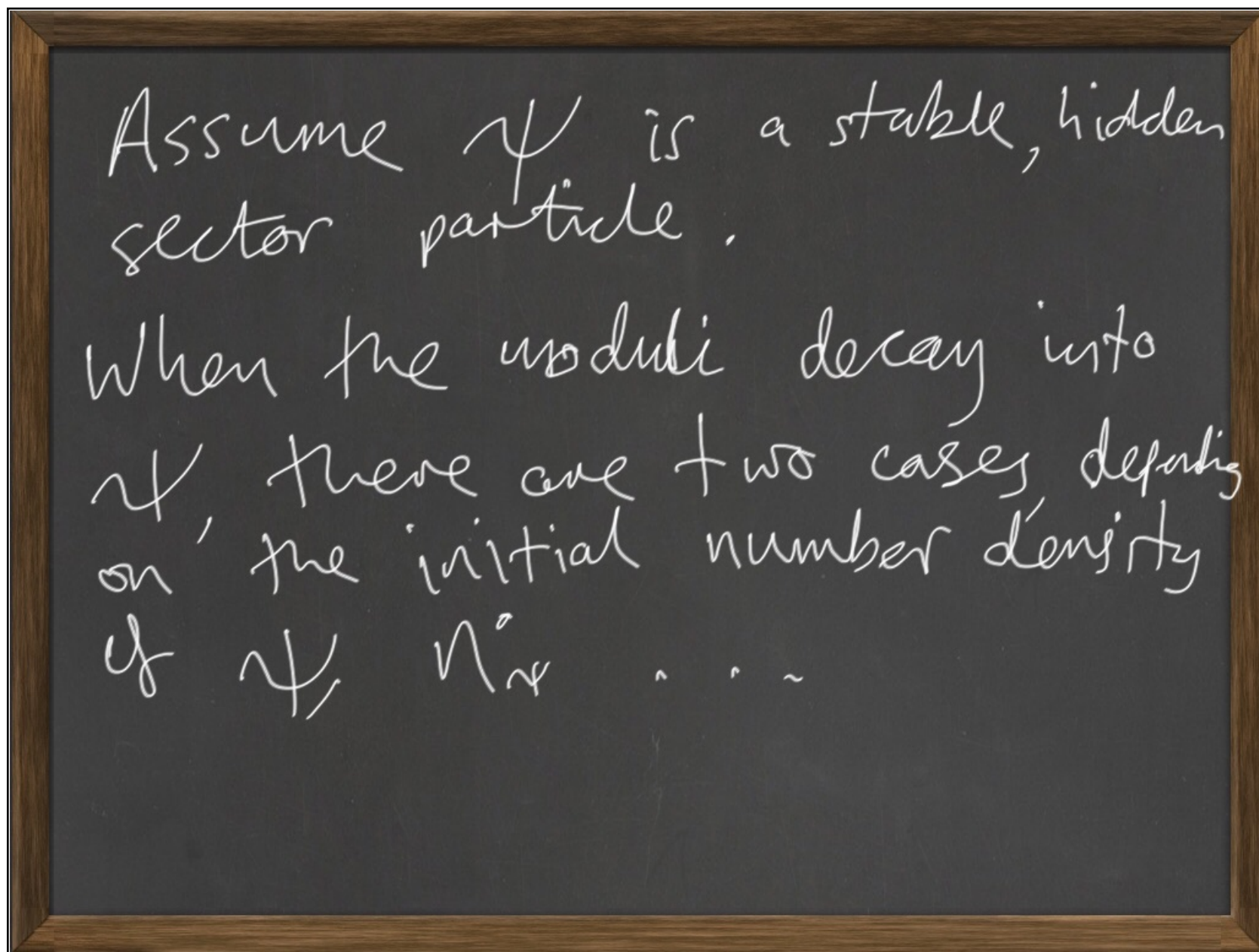
and $\tau_X \sim 10^{-17} \text{ s} \left(\frac{10^{-3}}{\epsilon}\right)^2 \times \text{mixing angles}$
for on shell Z

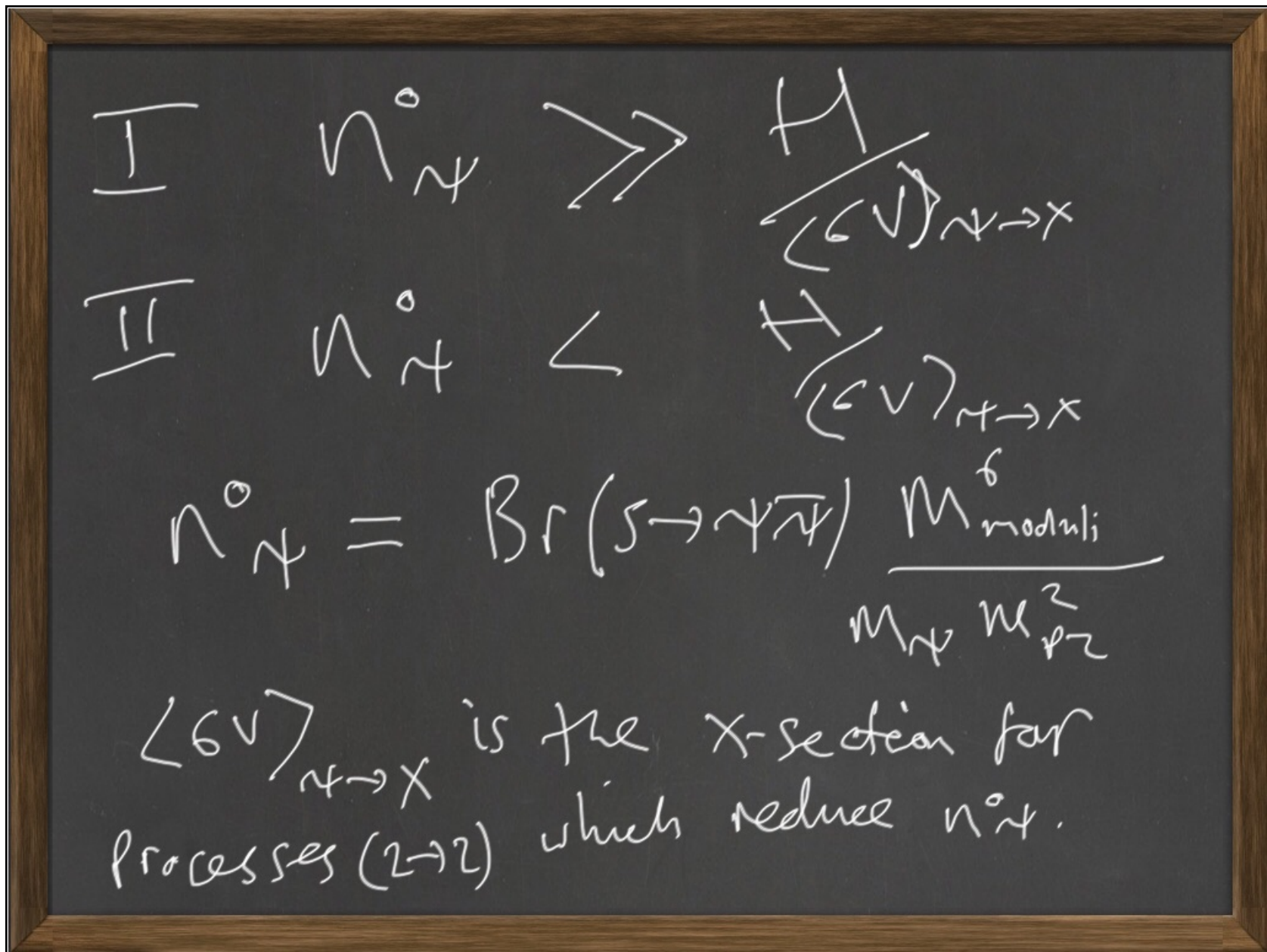
$\tau_X \sim \underline{10^{-9} \text{ s}} \left(\frac{10^{-3}}{\epsilon}\right)^2 \left(\frac{50 \text{ GeV}}{m_X - m_{X'}}\right)^4 \times \text{angles}$
for 3-body decay











Case I: ψ particles annihilate
 until $n_\psi = 3H / \langle \sigma v \rangle$.
 (work in progress)

Case II: ψ particles just hang
 around.

In case II $m_\psi \leq O(10^9) \text{ MeV}$.
 Heavier ψ ; give too much
 $\gg M$.

