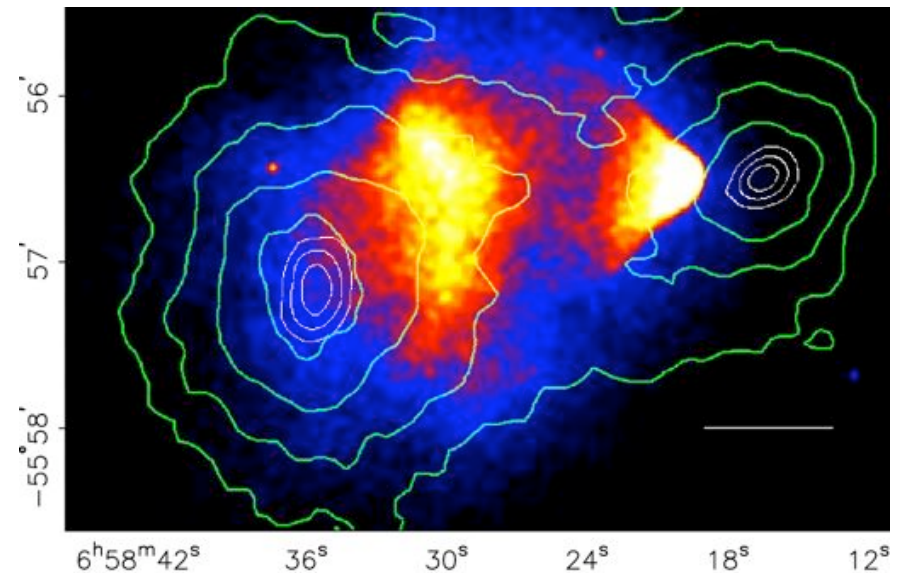
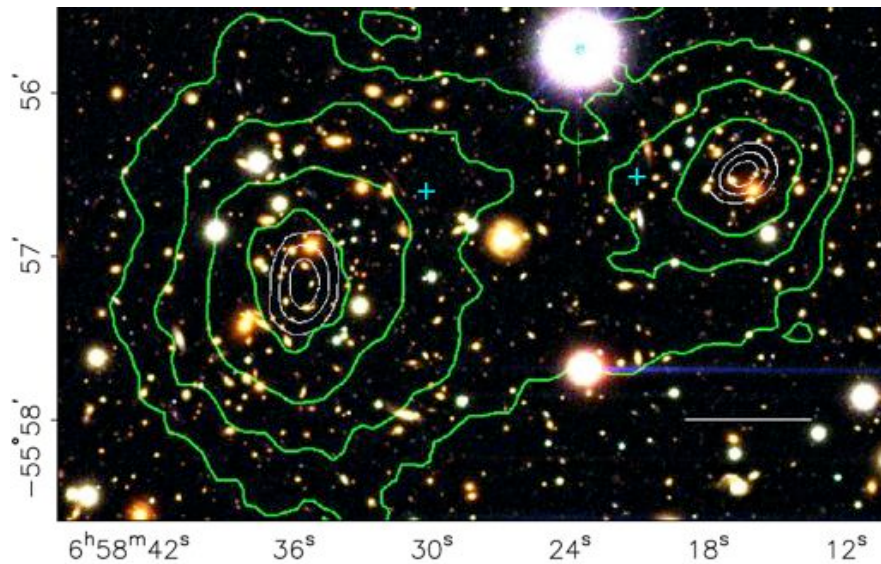


Modified Newtonian Dynamics as an alternative to dark matter: the Bullet Cluster case

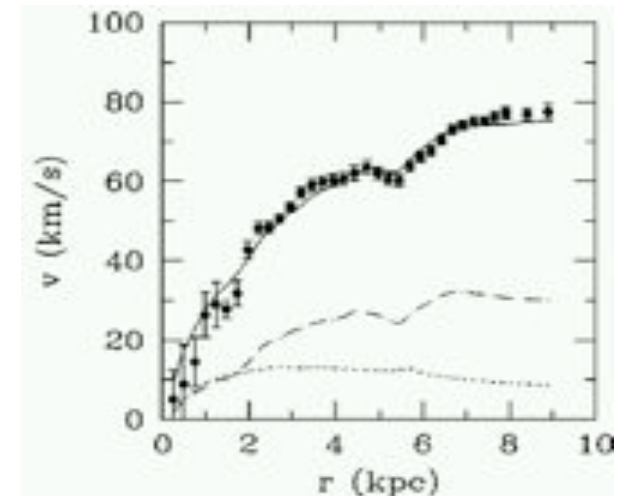
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In collaboration with G.W. Angus, H.S. Zhao (St Andrews), and H.Y. Shan (Beijing)



CDM: the cusp problem and the « correlation » problem

- Simulations of clustering CDM halos (e.g. Diemand et al.) predict a central cusp $\rho \propto r^{-\gamma}$, with $\gamma > 1$, observed neither in the MW (e.g. Famaey & Binney 2005), neither in HSB nor in LSB (No present-day satisfactory solution)
- Baryonic Tully-Fisher relation
 $V_{\infty}^4 \propto M_{\text{bar}}$ (tight->triaxiality of halo?)
- Tidal Dwarf Galaxies with DM?
(Gentile et al. [arXiv:0706.1976](#))
- **What is more:** wiggles of rotation curves follow wiggles of baryons in many HSB and in some LSB



Modified Newtonian Dynamics

- Correlation summarized by the **MOND** formula in galaxies (Milgrom 1983) :

$$\mu(|g|/a_0) \mathbf{g} = \mathbf{g}_{N \text{ baryons}} \quad \text{where } a_0 \sim cH_0$$

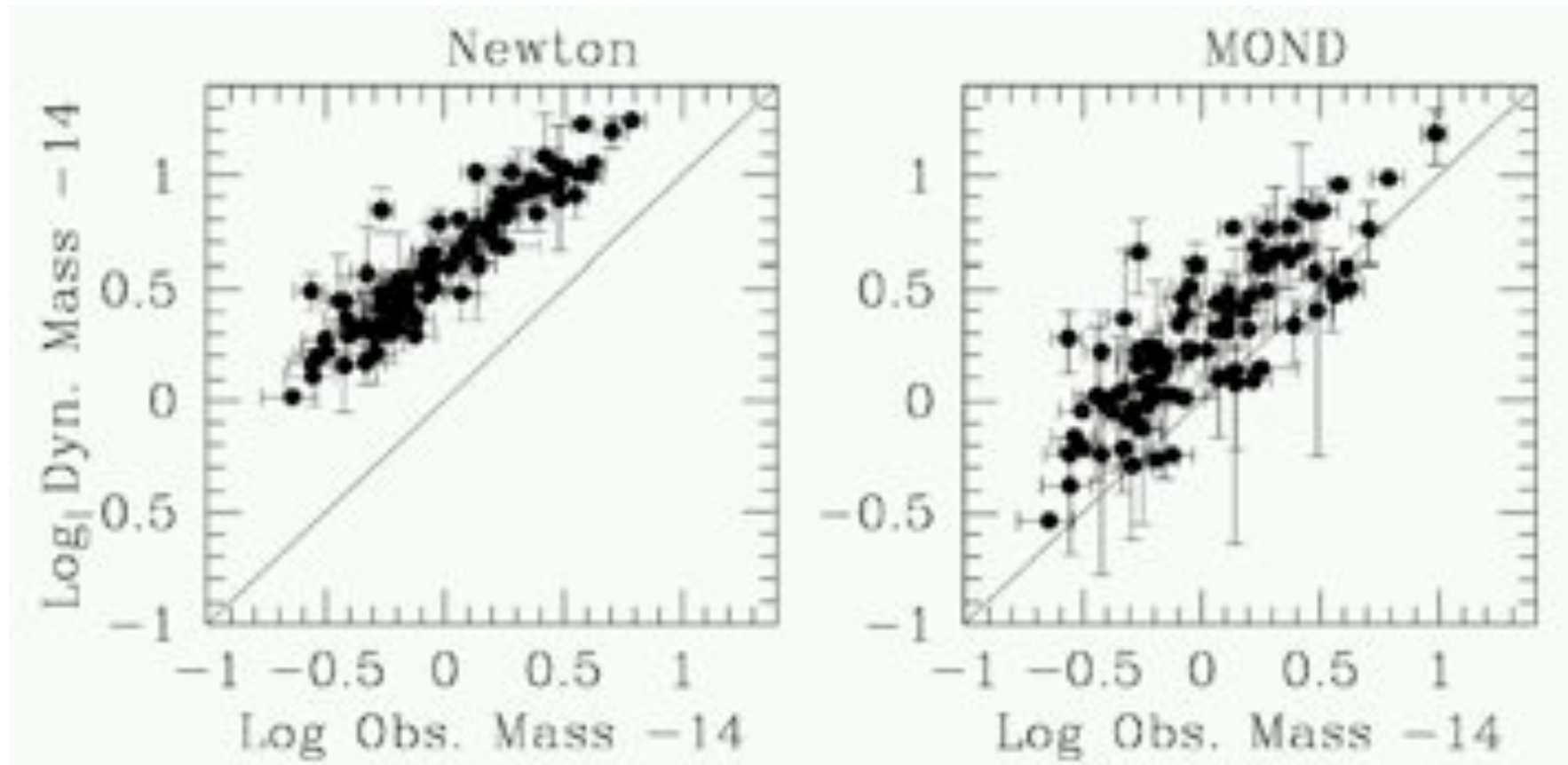
with $\mu(x) = x$ for $x \ll 1$ (MONDian regime) $\Rightarrow V^2/r \sim 1/r \Rightarrow V \sim \text{cst}$

$\mu(x) = 1$ for $x \gg 1$ (Newtonian regime)

- Why does it work in CDM and CDM-free galaxies?
- If fundamental: a) fundamental property DM ?
b) modification of gravity ?

$$\nabla \cdot [\mu (| \nabla \Phi | / a_0) \nabla \Phi] = 4 \pi G \rho$$

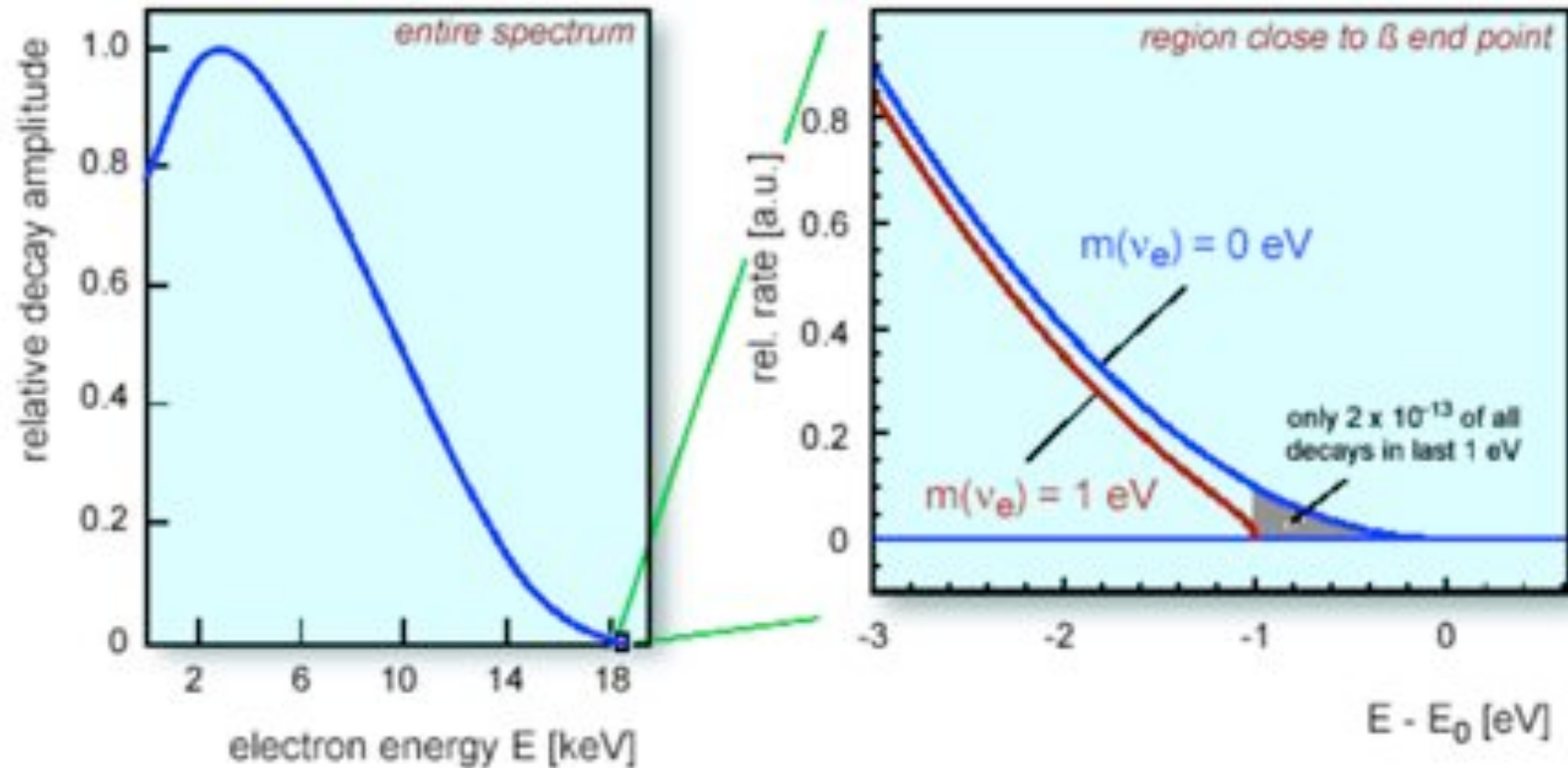
- Modifying GR to obtain **MOND** in **static weak-field limit**: dynamical 4-vector field $U^\alpha U_\alpha = -1$, with free function in the action playing the role of μ
(Bekenstein 2004; Zlosnik, Ferreira & Starkman 2007)



- But... rich clusters of galaxies need dark matter, e.g. ordinary neutrinos of 2eV (Sanders 2003, 2007) and/or dark baryons... 2eV neutrinos also invoked to fit the CMB power spectrum in order not to change the angular-distance relation (Skordis et al. 2006)
 Interestingly $\sum m_\nu \approx 6\text{eV}$ ($\Omega_\nu \approx 0.12$) excluded in ΛCDM

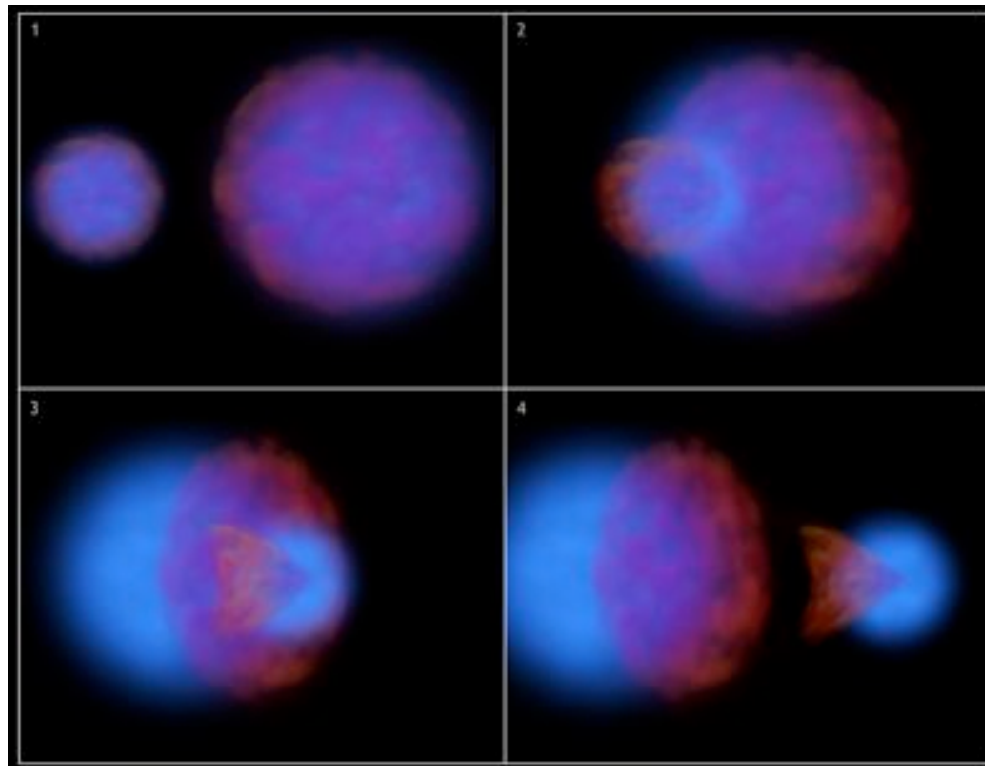
Mass of electron neutrino : KATRIN

β -decay of tritium (${}^3\text{H}$) into Helium 3 ion + electron + neutrino:



Which *new* challenge does the bullet cluster pose to MOND?

Merging galaxy cluster (Clowe et al., Bradač et al. 2006) with gas shock speed of 4500 km/s (mass centroids at smaller relative speed): a gigantic lab (1.4 Mpc for main axis) at a distance of 1Gpc ($z=0.3$), **separating the collisionless matter from the gas** (10^{13} and $2 \times 10^{13} M_{\text{sun}}$ of gas in the two cluster cores)



- We consider the weak-field static limit i.e. we neglect possible contributions of the spatial part of the vector field due to motion of matter sources (Dodelson & Liguori 2006)
- In the weak-field static limit there is a **linear** chain
 $\Phi \rightarrow g \rightarrow \alpha \rightarrow \kappa$
- In GR, there is an additional linear relation $\rho \rightarrow \Phi$ so the convergence $\kappa(R)$ directly measures the projected surface density $\Sigma(R)$
- In non-linear **MOND**, κ can be non-zero where there is **no projected matter** (Angus, Famaey & Zhao 2006)

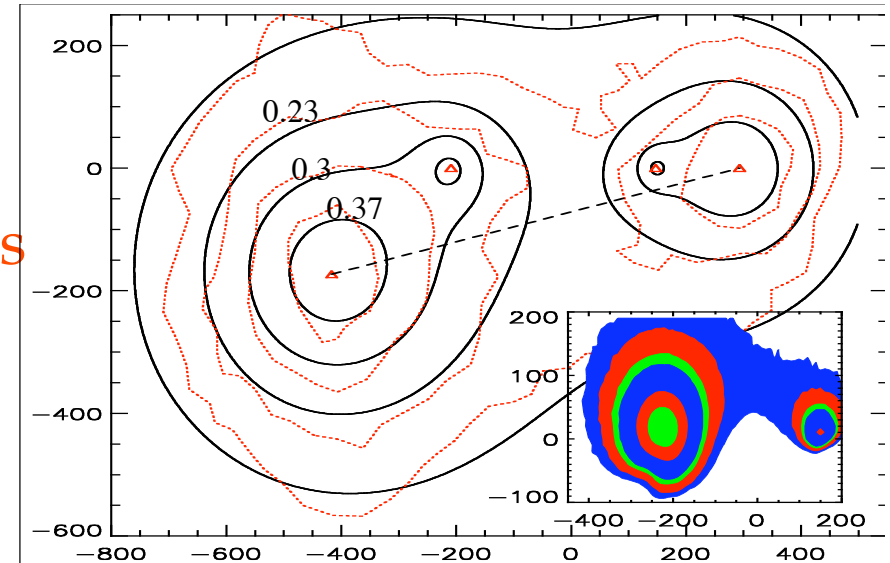
Angus, Shan, Zhao & Famaey (ApJ 654 L13)

- Take parametric logarithmic potential $\Phi(r)$

$$\Phi_i(r) = 1/2 v_i^2 \ln[1+(r/r_i)^2]$$

- Use $\Phi_1, \Phi_2, \Phi_3, \Phi_4$ for the 4 mass components of the bullet cluster

⇒ Parametric convergence $\kappa(R)$



- χ^2 fitting the 8 parameters on 233 points of the original weak-lensing convergence map

- With $\mu(x) = 1$ (\rightarrow GR), or e.g. $\mu(x) = x/(1+x)$, get enclosed $M(r)$:

$$4\pi GM(r) = \int \mu(|\nabla\Phi|/a_0) \partial\Phi/\partial r dA$$

- Collisionless: gas ratio within 180kpc of the galaxies and gas centers of the main cluster is **2.4:1** in **MOND**

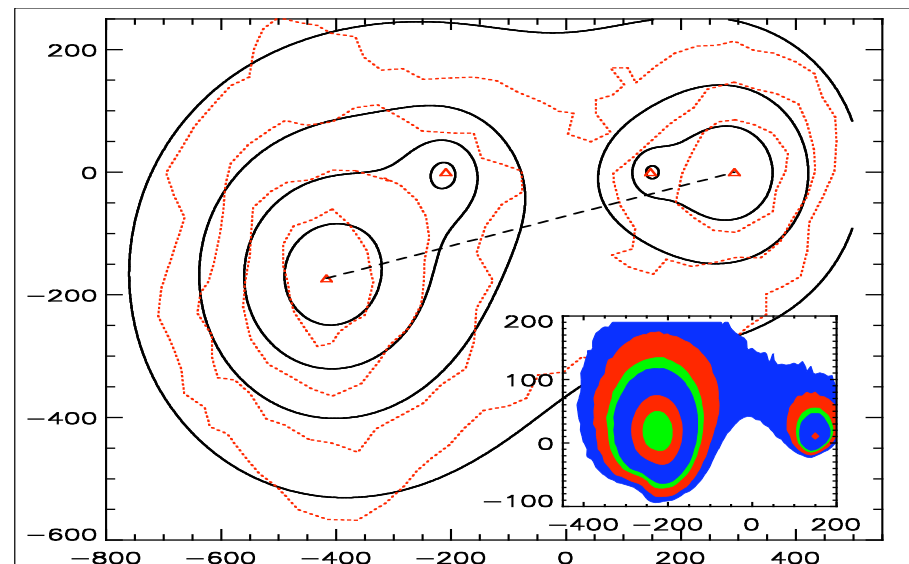
⇒ if **MOND** hidden mass = dark baryons

then they must be in **collisionless form** (e.g. dense clumps of cold gas)

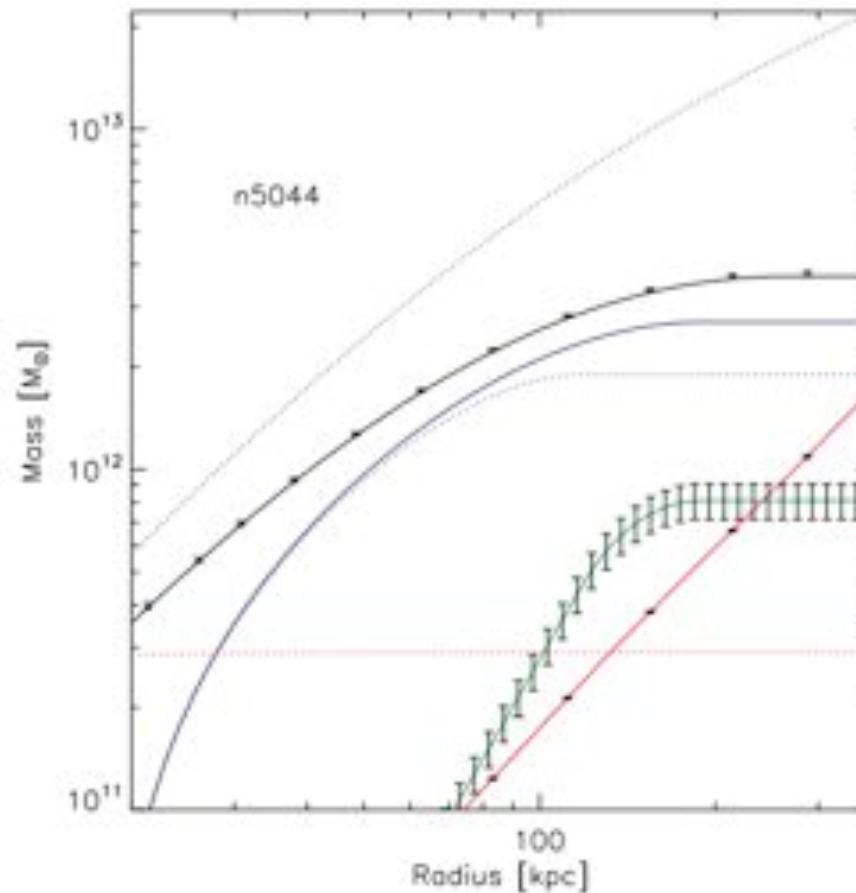
Missing baryons still represent $\sim 20\%$ of BBN (depending on the amount in WHIM), while baryonic mass in clusters only $\sim 10\%$ ⇒ room for dark baryons

- The **densities** of the collisionless matter in **MOND** are compatible with the maximum density of 2eV neutrinos! ($\sim 10^{-3} M_{\text{sun}}/\text{pc}^3$ in the bullet cluster for $T=9 \text{ keV} \sim 10^8 \text{ K}$)
⇒ not a new challenge to **MOND+neutrinos** (but problem from strong-lensing at center of subcluster)

- Problem for 2 eV neutrinos in a 100 kpc central core in rich clusters known from hydrostatic equilibrium since [Pointecouteau & Silk \(2005\)](#)
- Note that our potential-density model yields only the total mass => subtract the DM to get the gas BUT, of course, underestimate of the gas mass, especially for the subcluster



Angus, Famaey &
Buote 2007
arXiv:0709.0108



- Tremaine-Gunn limit for neutrinos:

$$\rho_{\nu} (\text{max}) \propto T^{3/2}$$

=> Problem for X-ray emitting groups with $T < 2$ keV

Conclusions

- The bullet cluster weak-lensing in itself is *not* a new challenge to **MOND + 2eV ordinary neutrinos**
- What about Abell 520? (Mahdavi et al. 2007) Effect of intercluster filaments on gravitational lensing in **MOND**?
- Looking forward to baby bullets' data
- Low-mass X-ray emitting groups are a tougher challenge:
 - => Maybe another fermionic dark particle?
(**hot sterile neutrinos** with $m_\nu \sim 10\text{eV}$?)
 - => Maybe **dark baryons** in the form of e.g. dense clumps of cold gas, present only in galaxy clusters? (but then, microlensing?)
 - => Maybe **CDM**?... 😊 but then one must understand why it does reproduce so precisely the **MOND** relation for all galaxies...